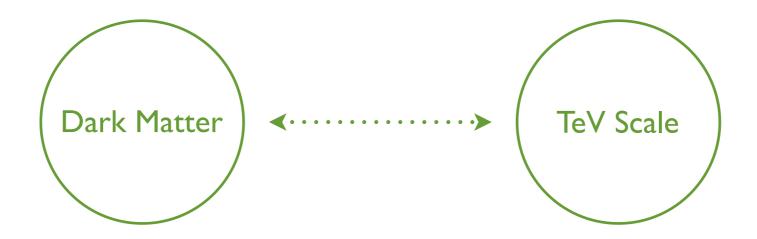
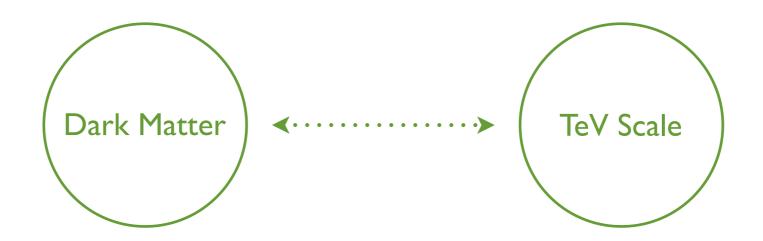
# Dark Matter Characterization at the LHC

Devin Walker
Harvard University

arXiv:0907.3142, arXiv:0907.3146, arXiv:1003.0899, and arXiv:1010.1??? with K.Agashe, D. Kim, M.Toharia and L. Zhu

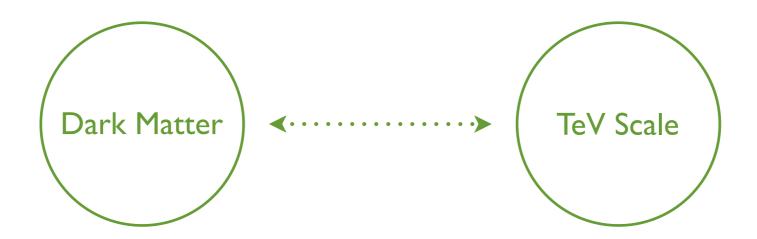


Connection with DM and TeV scale



- Connection with DM and TeV scale
  - Relic abundance for thermal DM

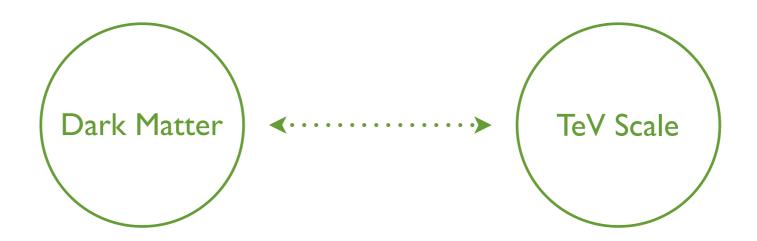
$$h^2 \Omega_{\rm DM} \simeq \frac{0.1 \ {\rm pb} \cdot c}{\langle \sigma v \rangle} \qquad \langle \sigma v \rangle \simeq \frac{g^4}{8\pi} \frac{1}{M^2}$$



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- Hierarchy Problem: TeV scale models often have O(100) GeV DM.
- LHC set to explore TeV scale!

 Broad Goal: Characterize LHC signatures with large missing energy to probe hidden sector/DM candidates.

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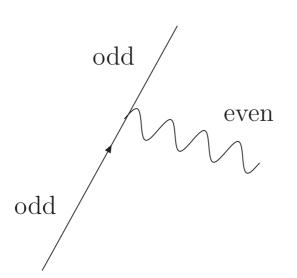
- Implement: Create new observables that can help theorists in reconstructing the model for dark matter.
- Today: Focus on distinguishing models with Z2 (parity) stabilization symmetries from models with other stabilization symmetries.

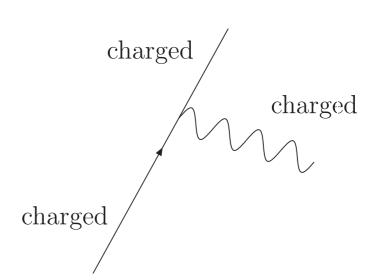


Ultimately searching for different decay topologies

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  - Main point: Different stabilization symmetries generate different decay topologies.

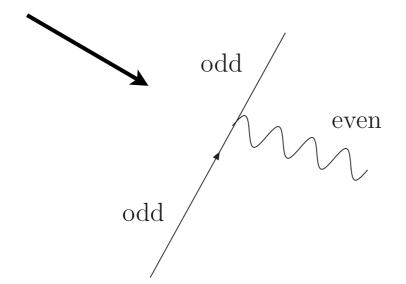
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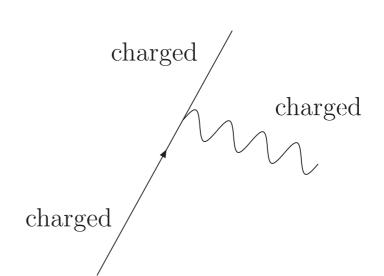




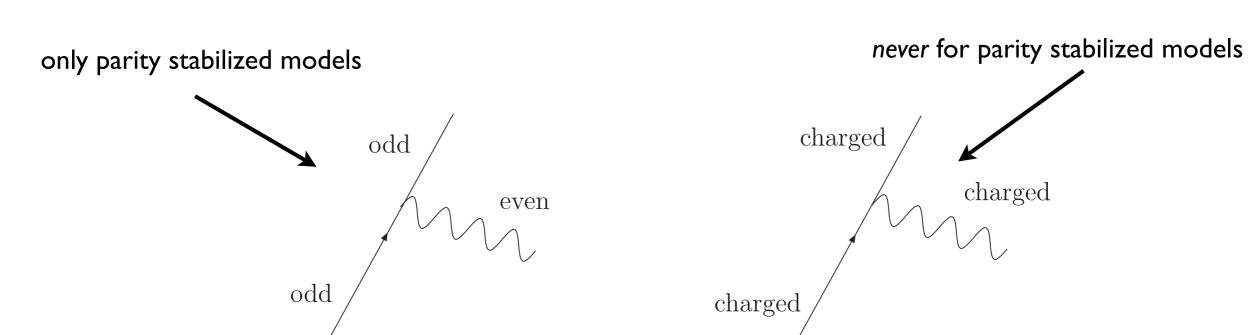
Ultimately searching for different decay topologies

only parity stabilized models





Ultimately searching for different decay topologies



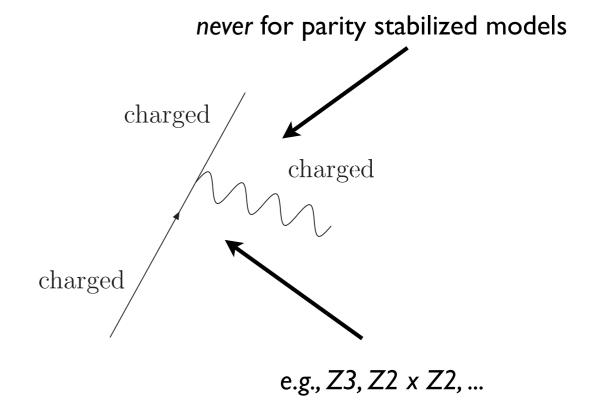
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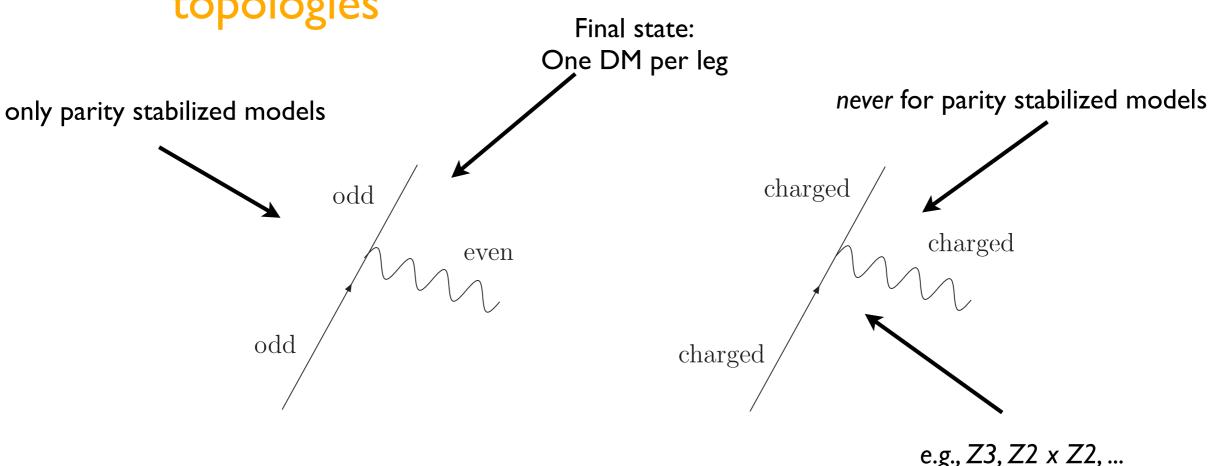
odd

even

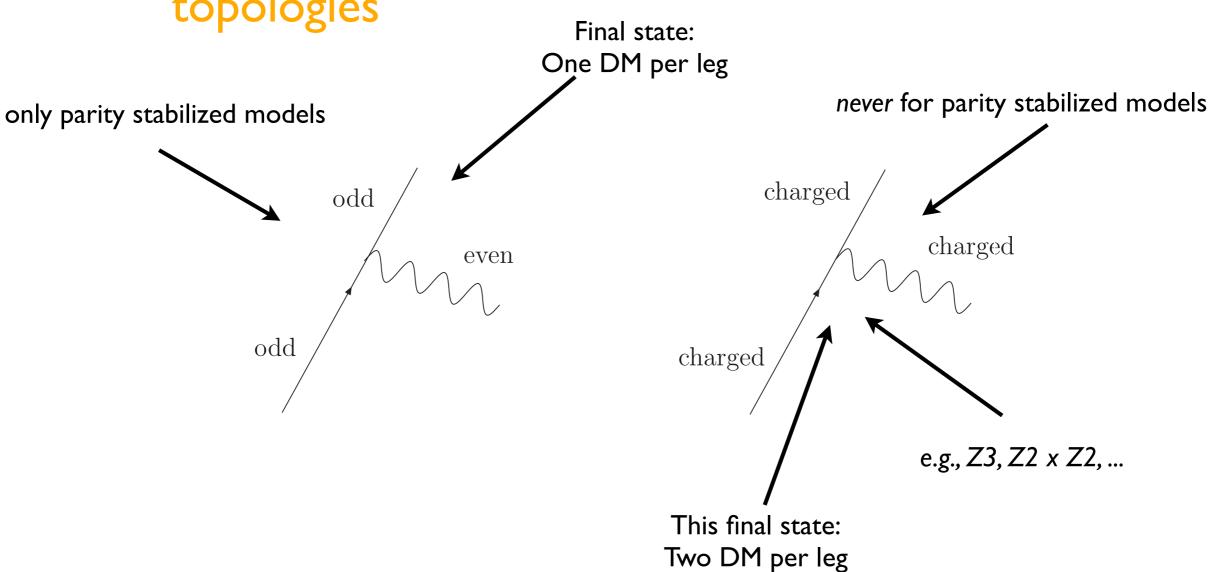
odd



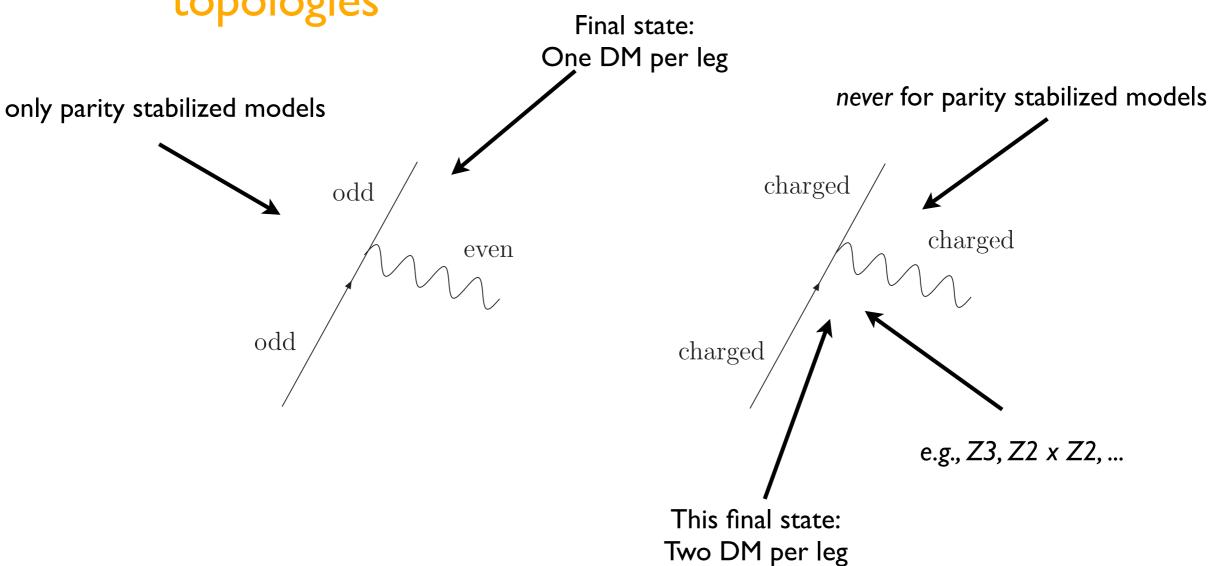
Ultimately searching for different decay topologies



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Ultimately searching for different decay topologies



Searching for scenarios w/ one or two DM per leg.

 Unique Decay Topology and Reconstruction (10 Minutes)

 Unique Decay Topology and Reconstruction (10 Minutes)

 Topologies and MT2 Reconstruction (10 Minutes)

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 Signatures with Metastable Particles (7 Minutes)

 Unique Decay Topology and Reconstruction (10 Minutes)

 Topologies and MT2 Reconstruction (10 Minutes)

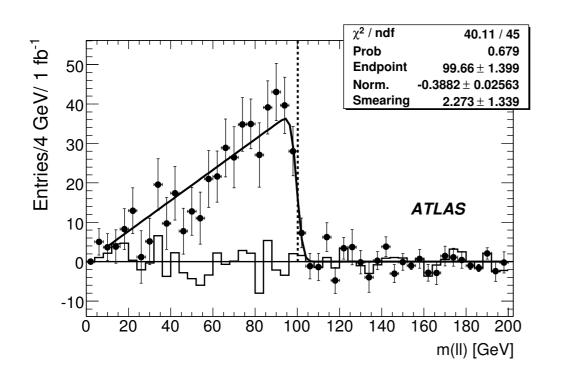
 Signatures with Metastable Particles (7 Minutes)

Future work/Conclusions

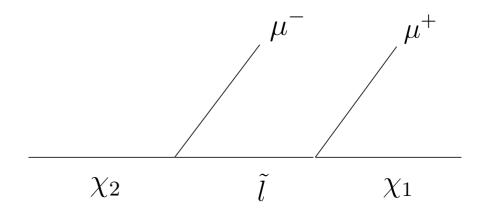
# Part I: Unique Decay Topologies and Reconstruction

## Decays for Parity Models (a one slide review)

Standard searches at LHC involves kinematic endpoints.



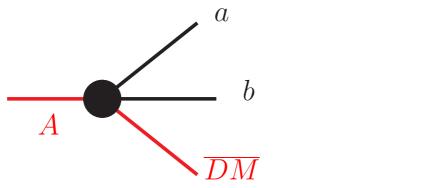
$$\mathrm{m}_{\chi_2}$$
 = 218.6 GeV,  $\mathrm{m}_{\tilde{l}}$  = 230.45,  $\mathrm{m}_{\chi_1}$  = 117.91

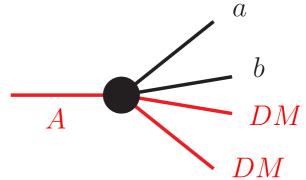


$$m_{\ell\ell}^{\text{edge}} = m_{\tilde{\chi}_2^0} \sqrt{1 - \left(\frac{m_{\tilde{\ell}}}{m_{\tilde{\chi}_2^0}}\right)^2} \sqrt{1 - \left(\frac{m_{\tilde{\chi}_1^0}}{m_{\tilde{\ell}}}\right)^2}$$

## Non-Parity Decay Chains

"Non-parity" models have one or more DM candidates per decay chain!





 $\mathcal{Z}_2$  and non-parity diagrams

A correction for "non-parity" diagrams

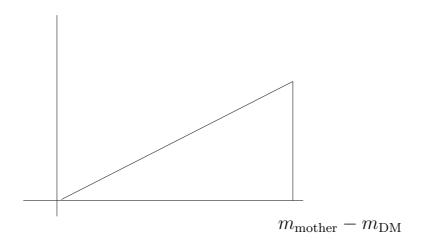
## Non-Parity Decay Chains

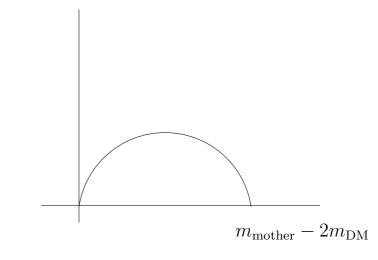
"Non-parity" models have one or more DM candidates per decay chain!

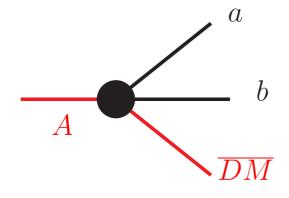


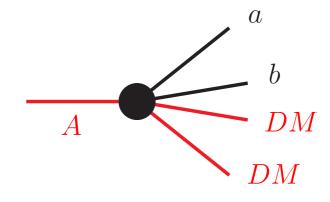
- Generates a double kinematic edge.
- Direct calculation of DM mass.

#### Off-shell cartoon:



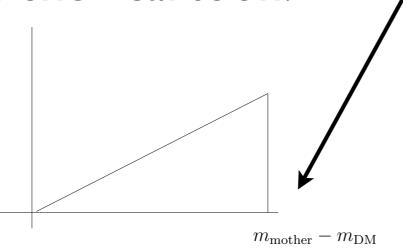


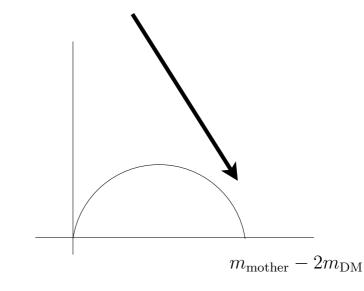


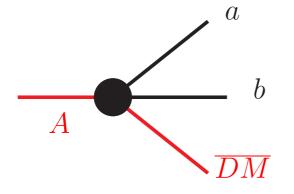


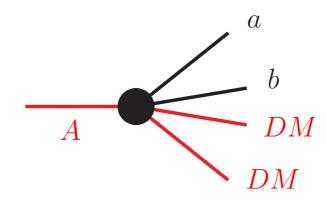
Different edges and shapes

Off-shell cartoon:

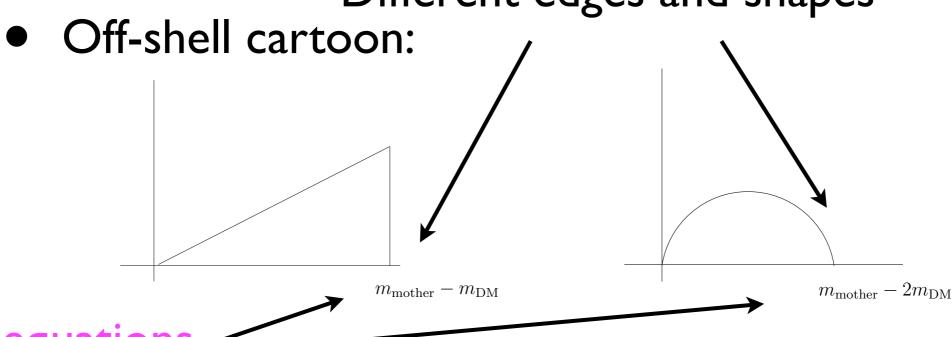








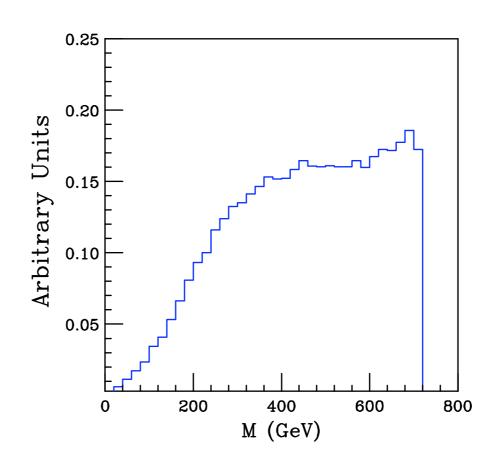
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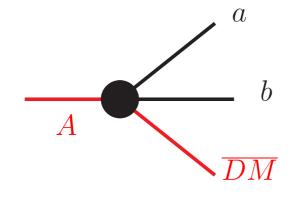


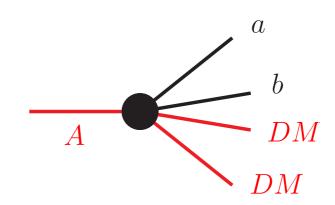
Two equations - two unknowns



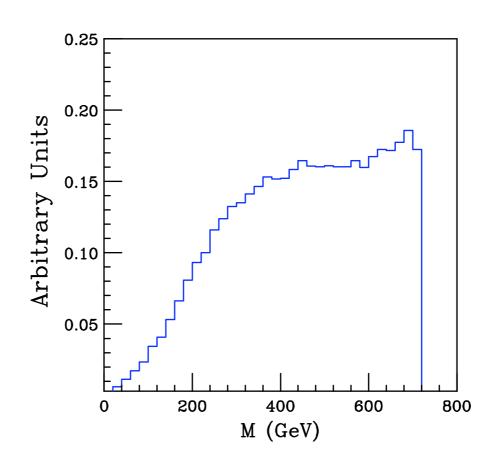
Cross sections summed





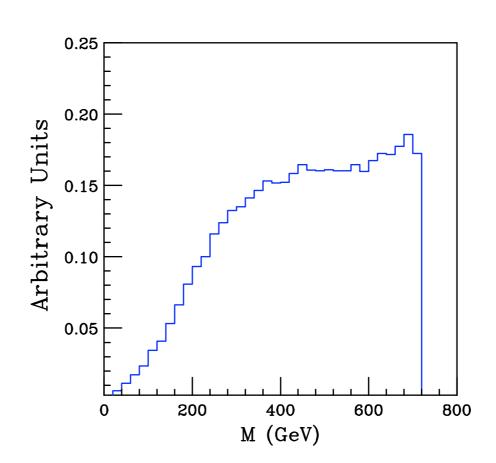


Cross sections summed



 820 GeV mother and 100 GeV DM, vector-like couplings.

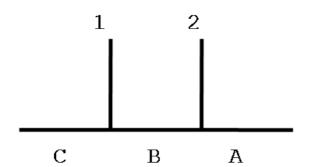
Cross sections summed

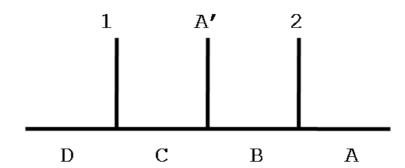


- 820 GeV mother and 100 GeV DM, vector-like couplings.
- Second endpoint at about 620 GeV.

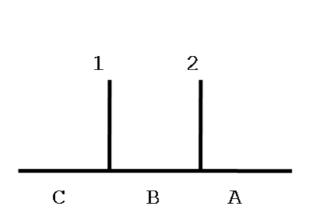
## On-Shell Considerations

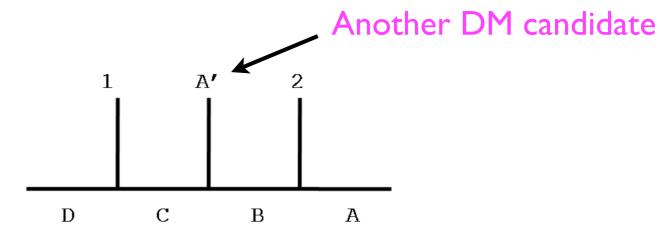
 Any unique features underneath the effective coupling?



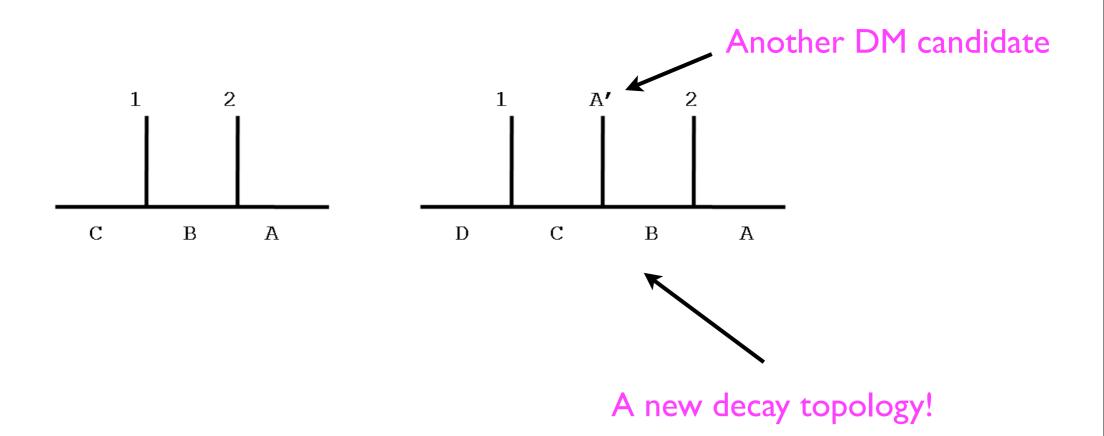


$$m_D > m_C > m_B > m_A$$

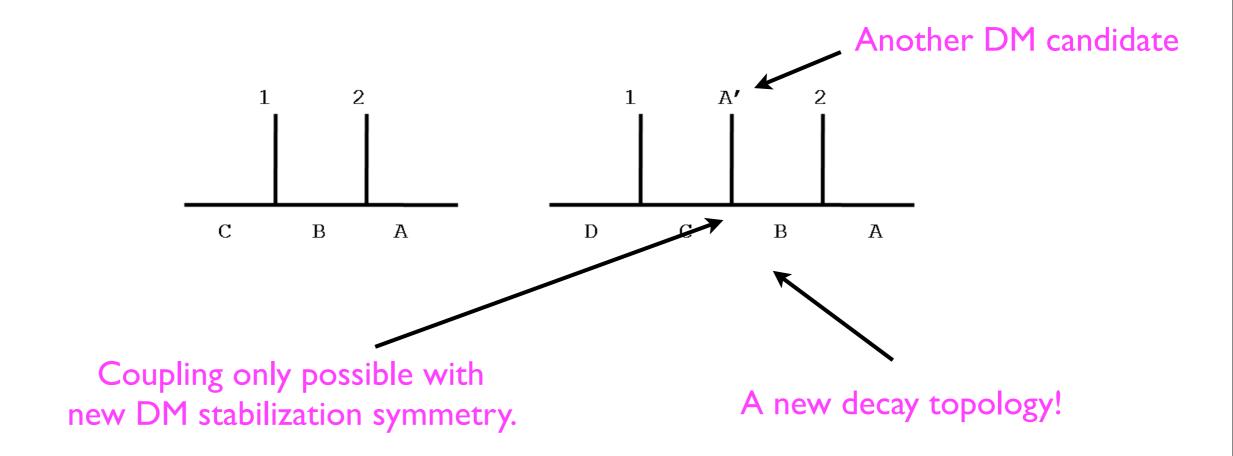




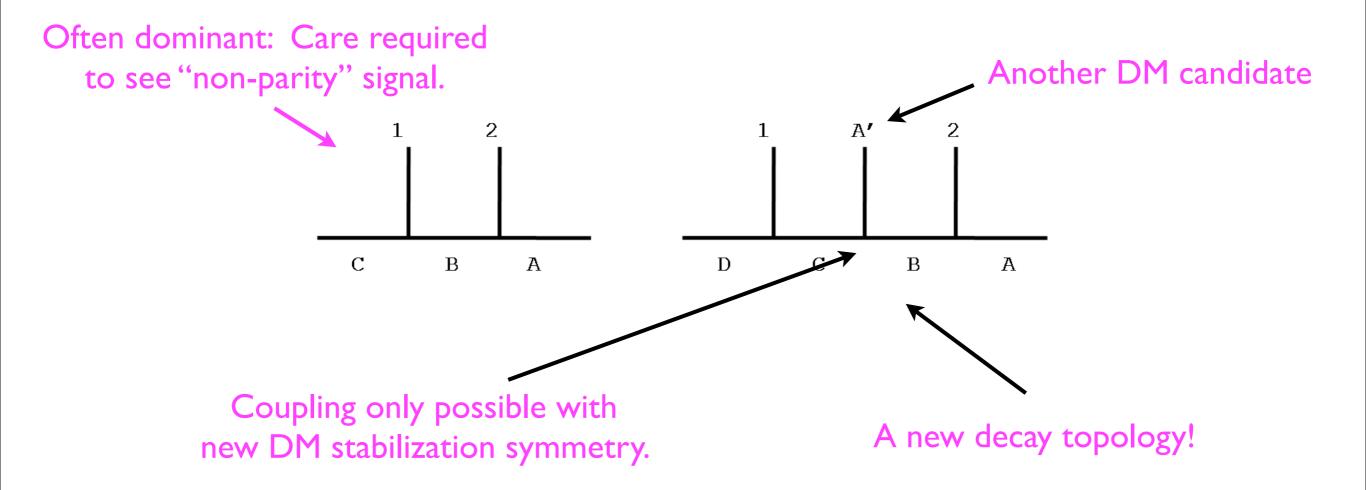
$$m_D > m_C > m_B > m_A$$



$$m_D > m_C > m_B > m_A$$



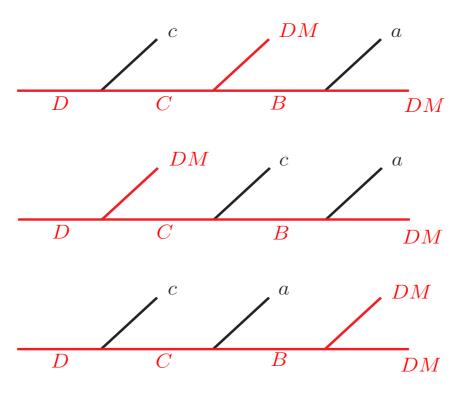
$$m_D > m_C > m_B > m_A$$



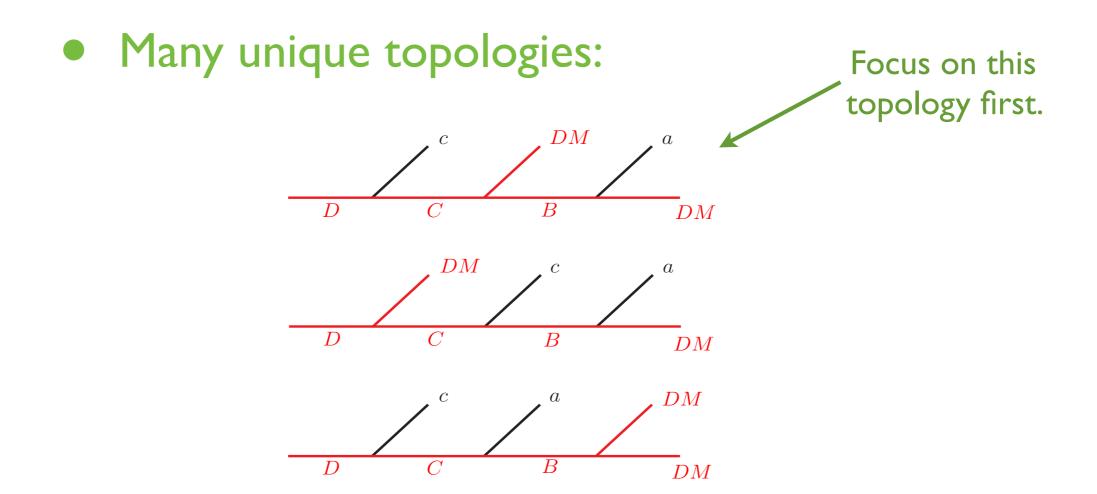
$$m_D > m_C > m_B > m_A$$

## Additional Possible Topologies

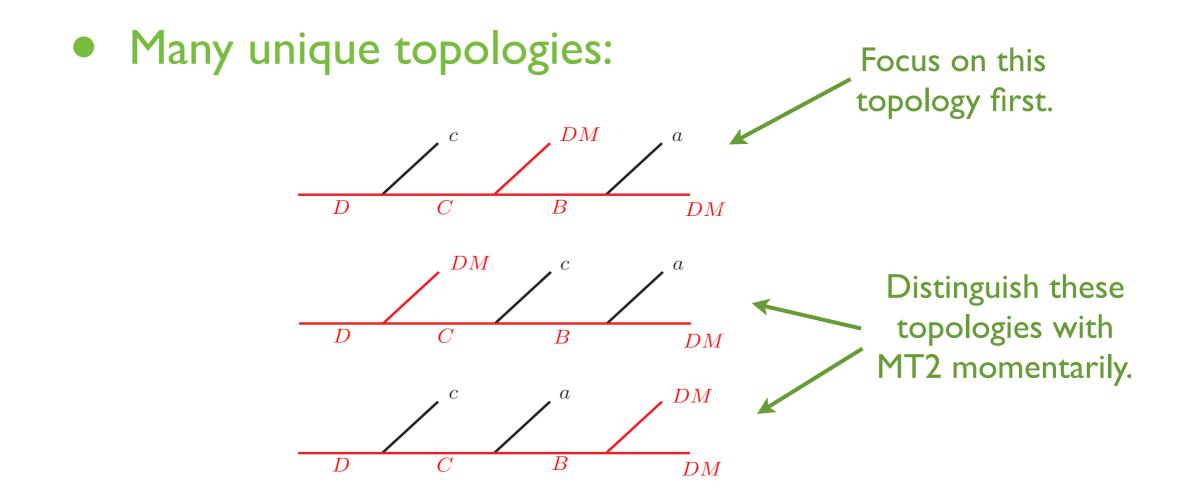
Many unique topologies:



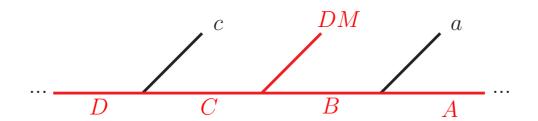
## Additional Possible Topologies

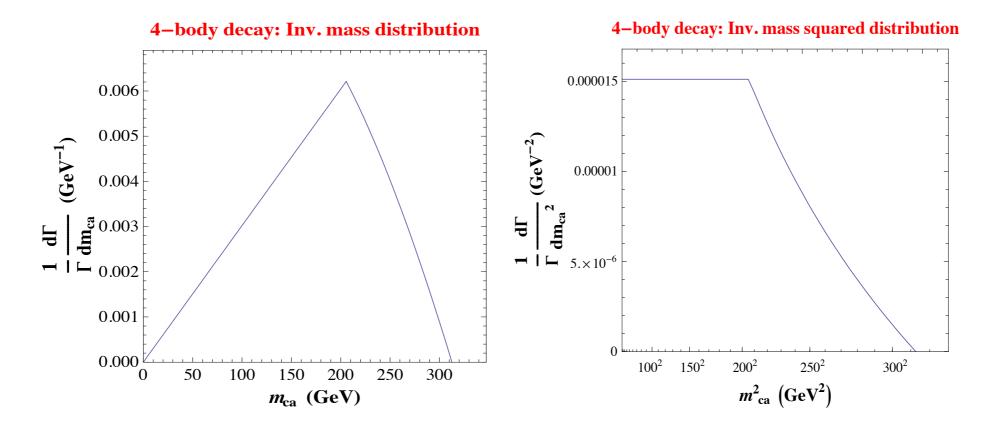


## Additional Possible Topologies



## Our Signal

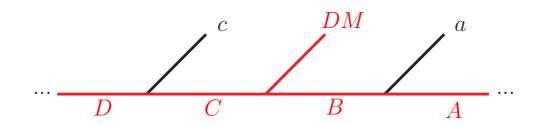


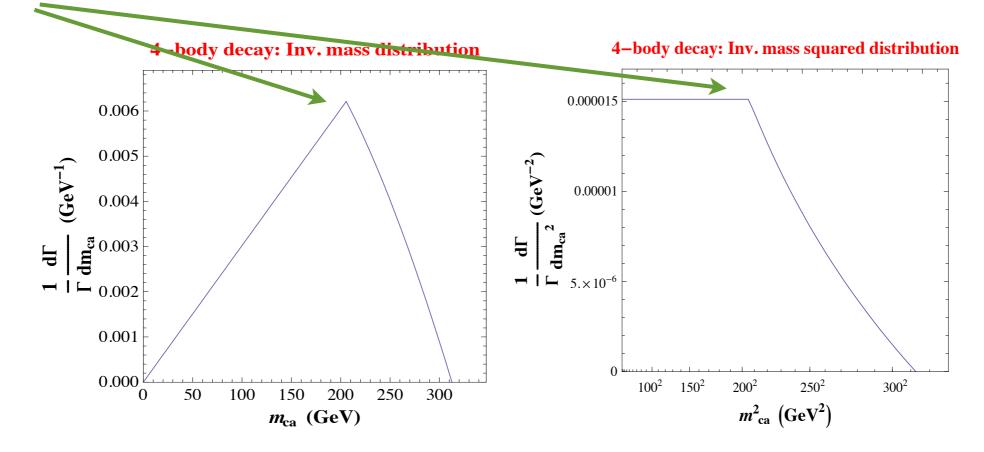


A, B, C and D have masses of 200, 400, 700 and 800 GeV

## Our Signal

Discontinuity in the cross section.

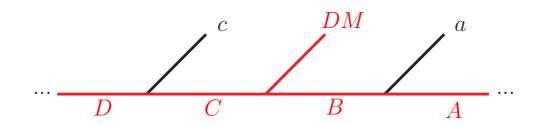


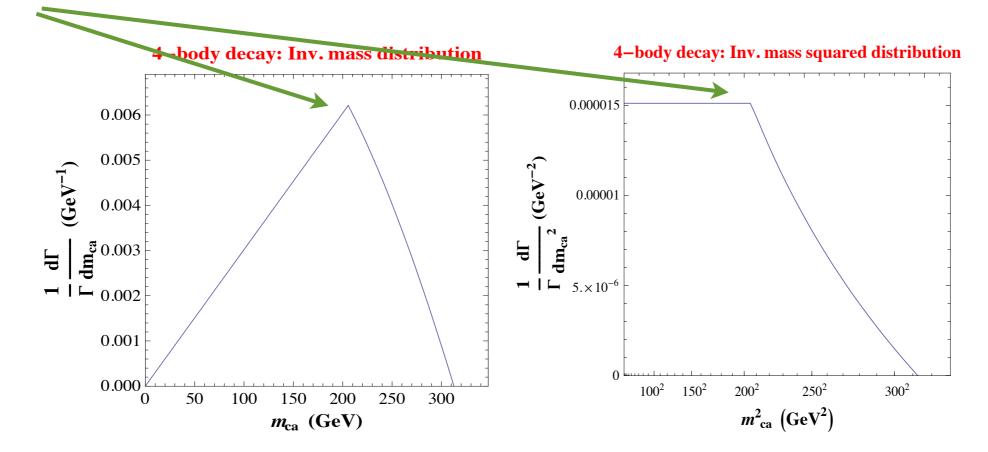


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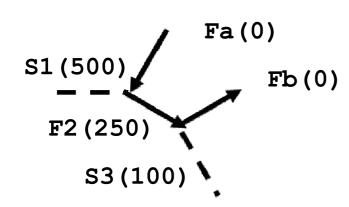




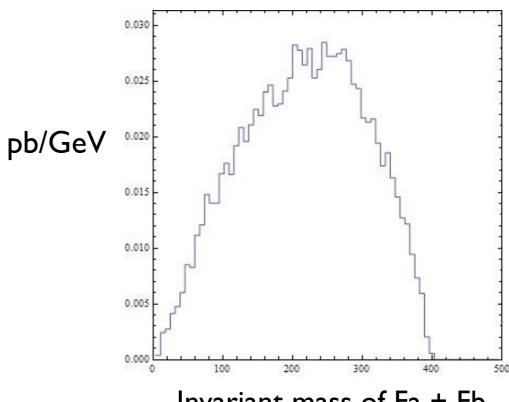
A, B, C and D have masses of 200, 400, 700 and 800 GeV

Spin correlations important for signal!

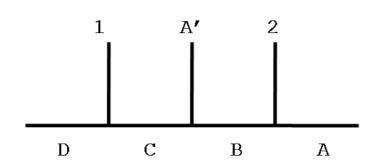
## The Importance of Spin

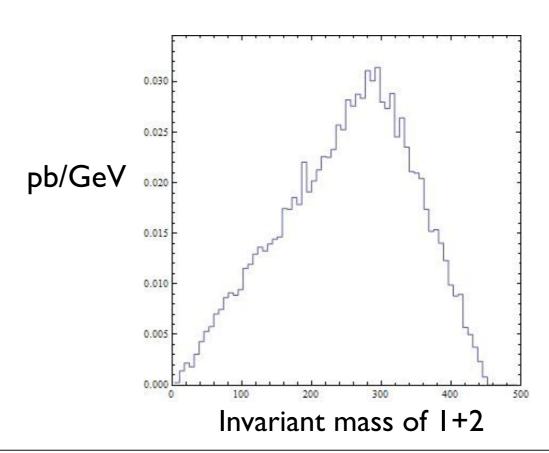


Left Chiral - Right Chiral

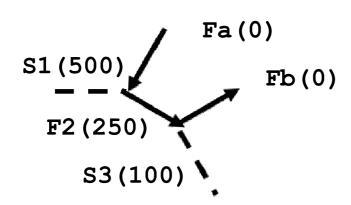


Invariant mass of Fa + Fb





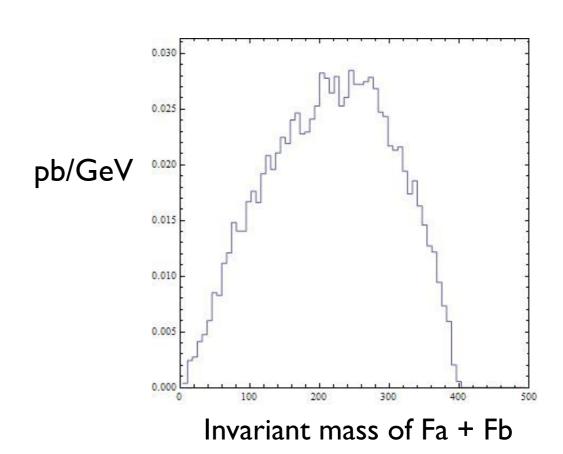
## The Importance of Spin

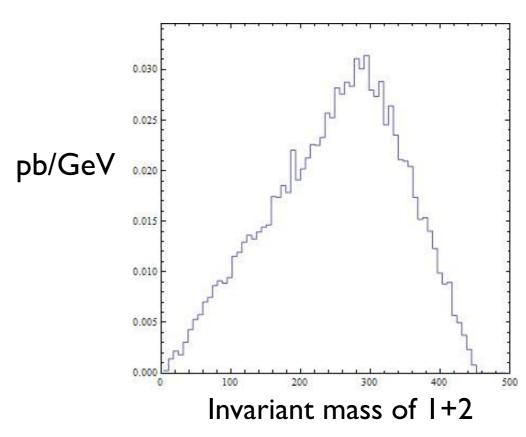


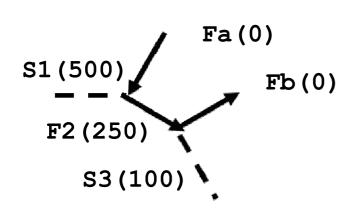
Left Chiral
- Right Chiral

Analytically there is a derivative discontinuity in the cross section...

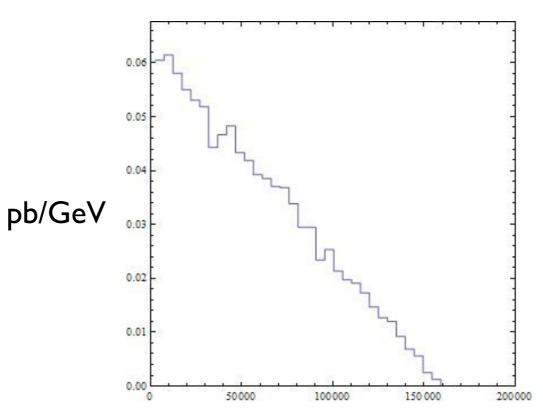
(no spin correlations)



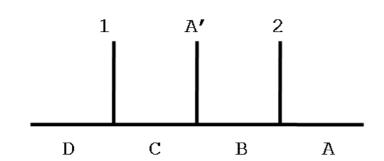


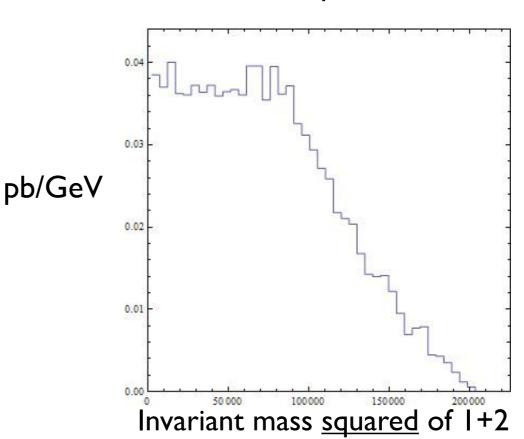


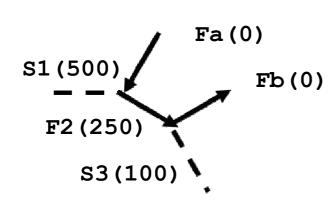
Left Chiral
- Right Chiral



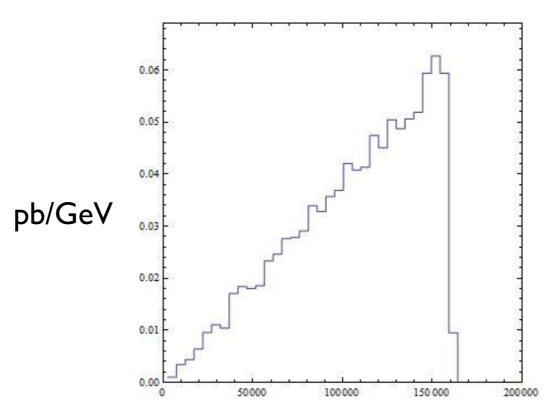
Invariant mass squared of Fa + Fb



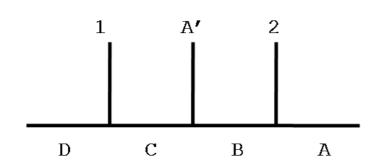


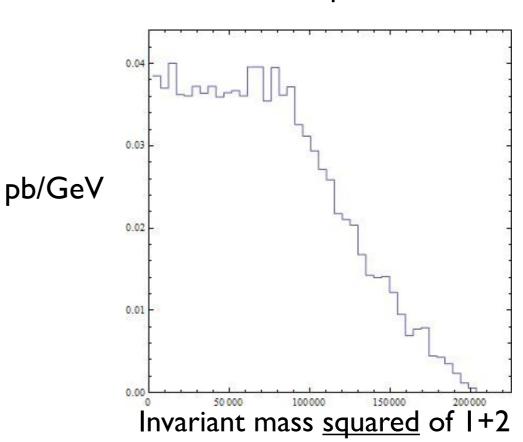


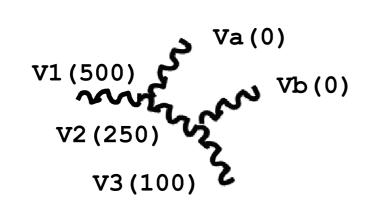
Left Chiral
- Left Chiral

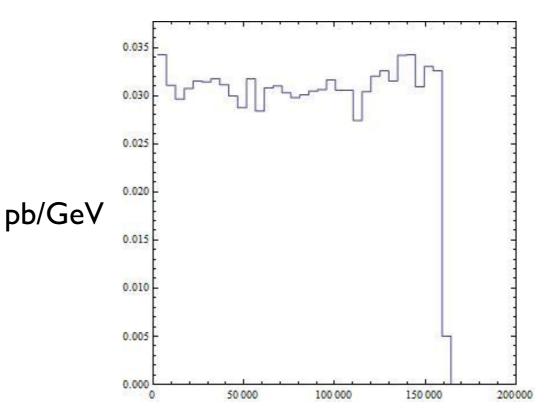


Invariant mass squared of Fa+Fb

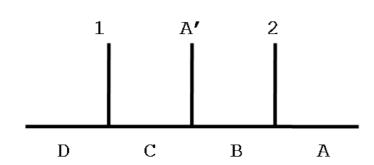


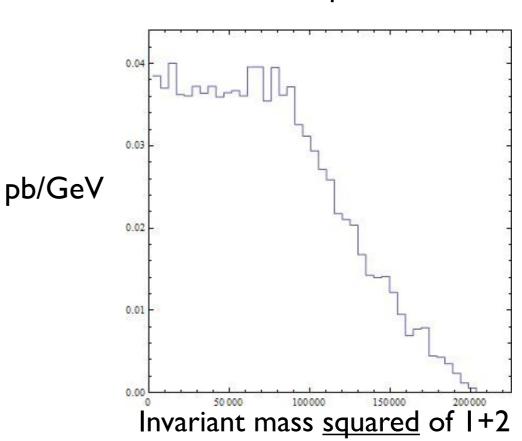


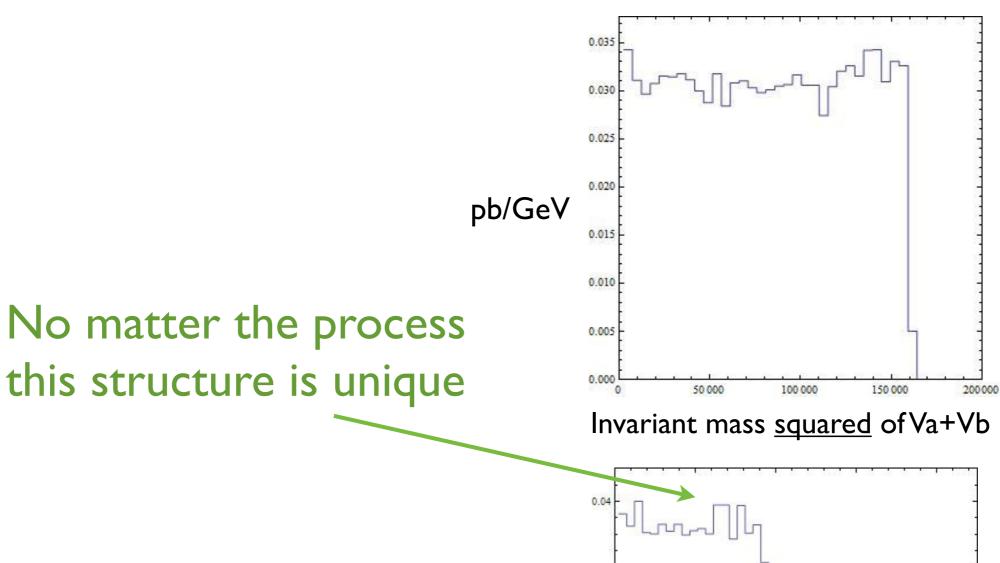


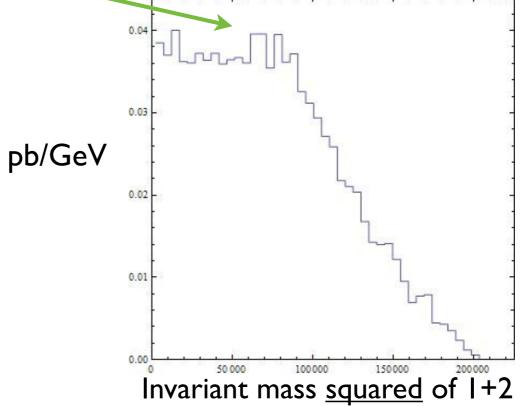


Invariant mass squared of Va+Vb

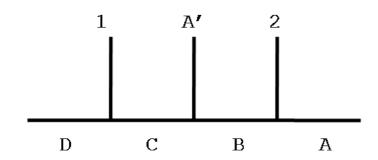


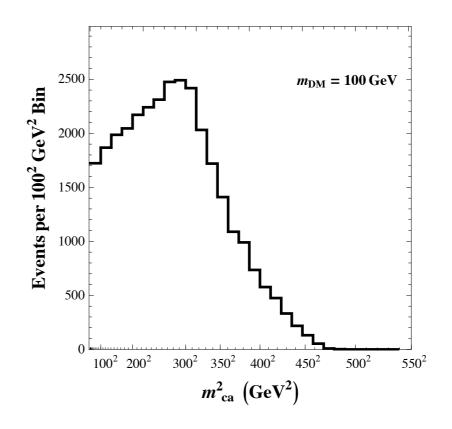






Our signal process (with spin correlations)

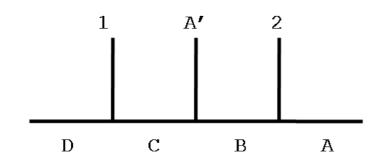


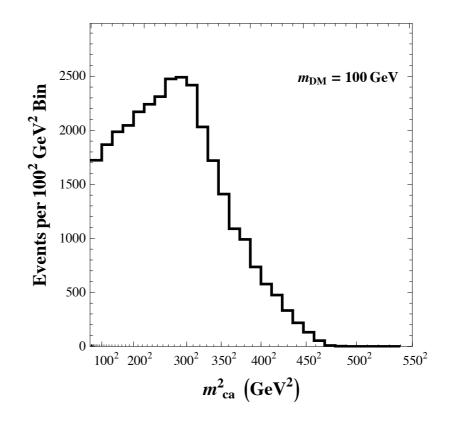


Invariant mass <u>squared</u> of I+2

Exotic example: B and C is spin I and I/2.
 Chiral couplings.

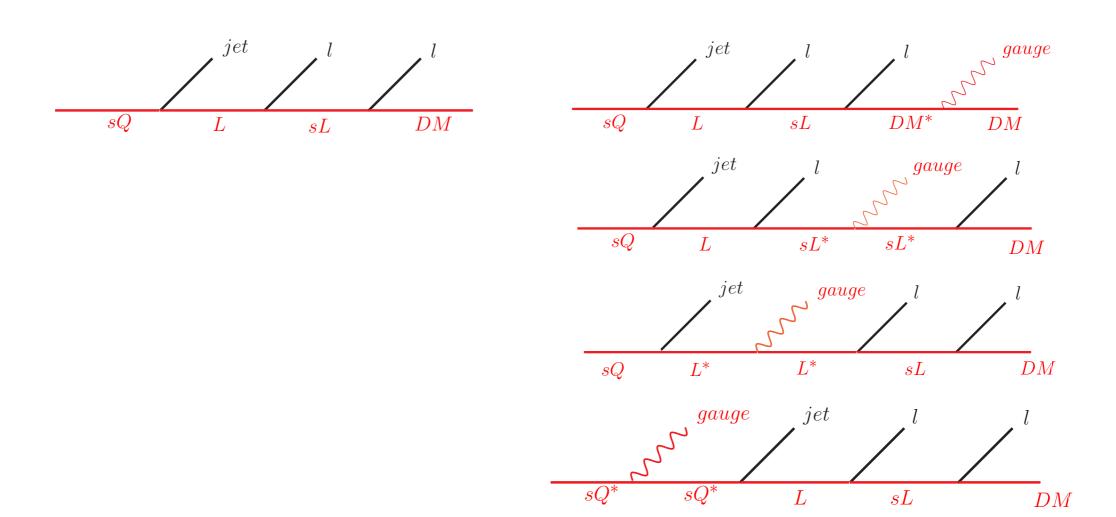
Our signal process (with spin correlations)





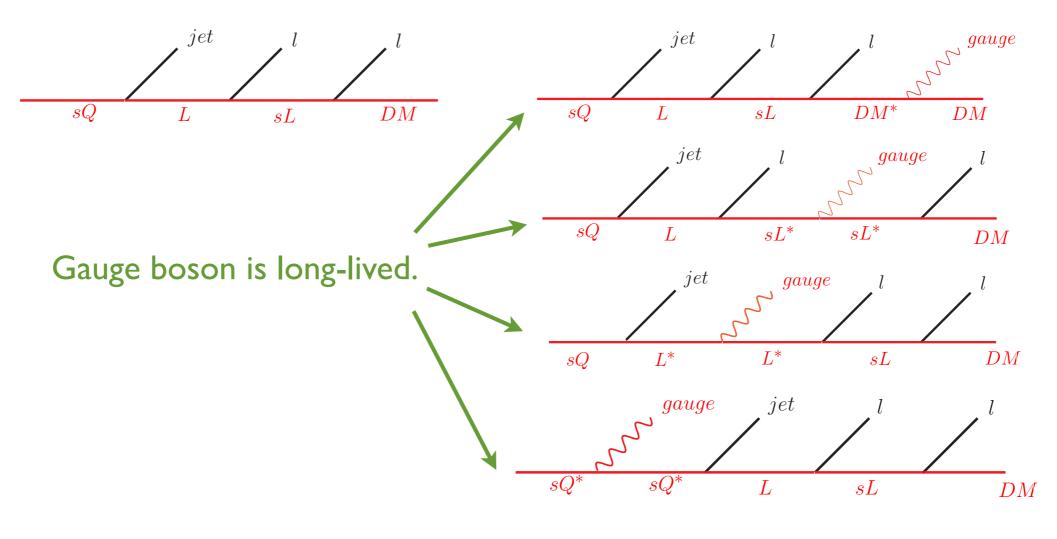
Invariant mass <u>squared</u> of I+2

- Exotic example: B and C is spin I and I/2.
   Chiral couplings.
- How does this feature hold up to acceptance and detector cuts?

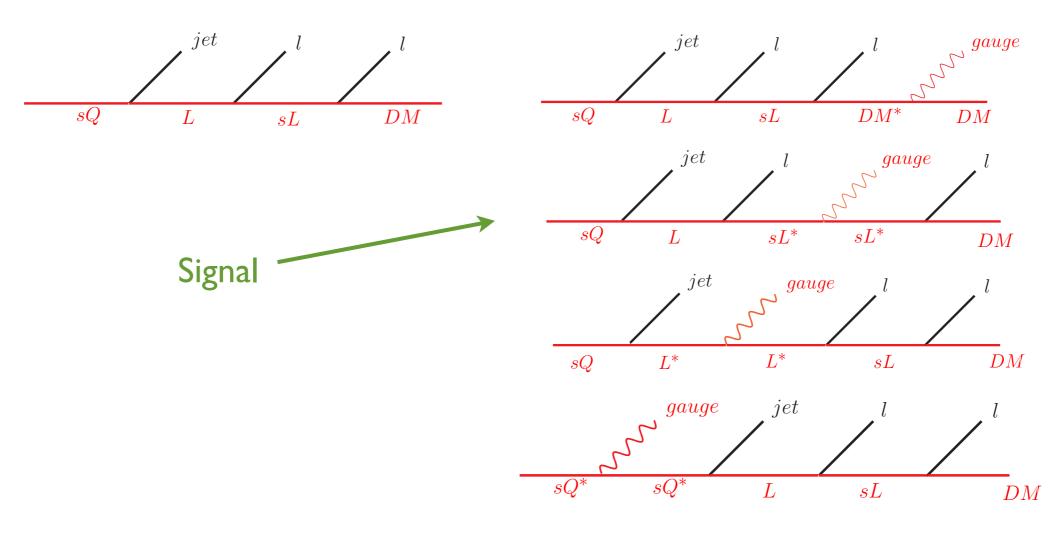


$$m_Q = 700 \; \text{GeV} \quad m_L = 650 \; \text{GeV} \quad m_{sL} = 300 \; \text{GeV} \quad m_\chi = 100 \; \text{GeV} \quad m_V = 100 \; \text{GeV}$$

Consider the following processes:



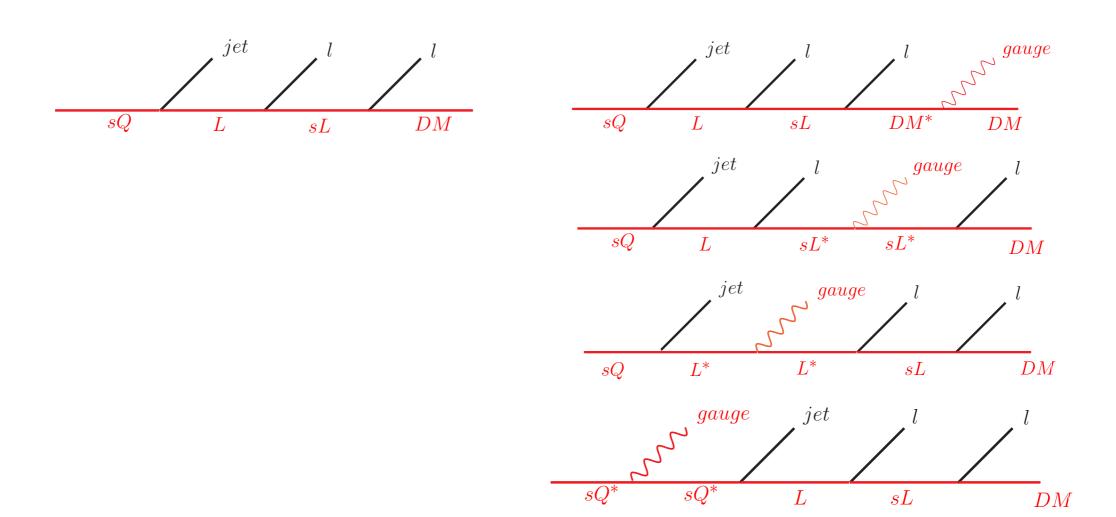
 $m_Q = 700 \; \text{GeV} \quad m_L = 650 \; \text{GeV} \quad m_{sL} = 300 \; \text{GeV} \quad m_\chi = 100 \; \text{GeV} \quad m_V = 100 \; \text{GeV}$ 



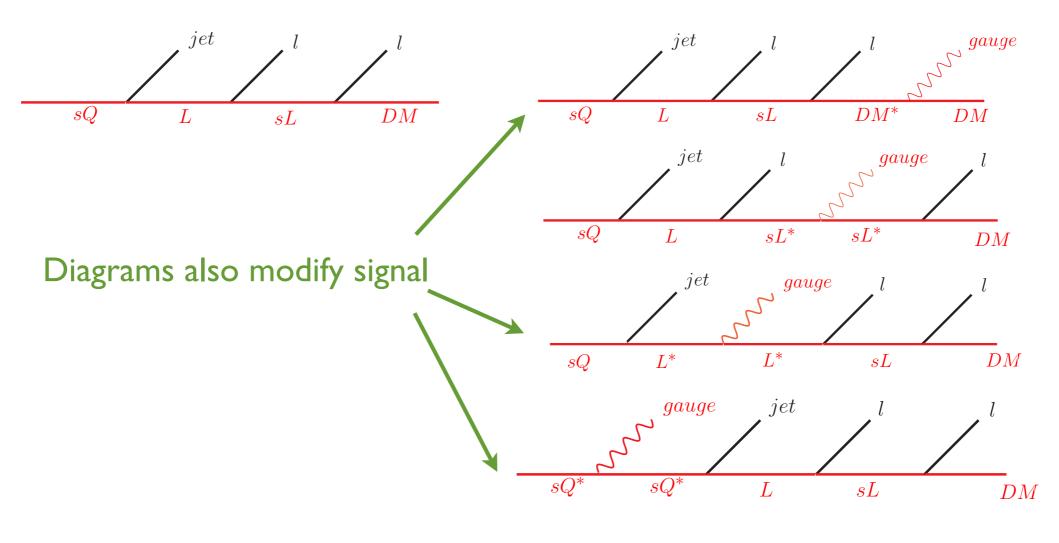
$$m_Q = 700 \text{ GeV}$$
  $m_L = 650 \text{ GeV}$   $m_{sL} = 300 \text{ GeV}$   $m_\chi = 100 \text{ GeV}$   $m_V = 100 \text{ GeV}$ 

$$m_{\chi} = 100 \text{ GeV}$$

$$m_V = 100 \text{ GeV}$$



$$m_Q = 700 \; \text{GeV} \quad m_L = 650 \; \text{GeV} \quad m_{sL} = 300 \; \text{GeV} \quad m_\chi = 100 \; \text{GeV} \quad m_V = 100 \; \text{GeV}$$



$$m_Q = 700 \text{ GeV}$$
  $m_L = 650 \text{ GeV}$   $m_{sL} = 300 \text{ GeV}$   $m_\chi = 100 \text{ GeV}$   $m_V = 100 \text{ GeV}$ 

$${
m GeV}$$

$$m_{sL} = 300 \text{ GeV}$$

$$m_{\chi} = 100 \text{ GeV}$$

$$m_V = 100 \text{ GeV}$$

 Acceptance and background cuts consistent with ATLAS and CMS:

$$|\eta_l| < 2.5,$$
  $|\eta_j| < 2.5,$   $\Delta R_{lj}, \Delta R_{jj} > 0.4.$ 

- 1. Two leptons with  $p_T > 20 \text{ GeV}$
- 2. At least one leading jet with  $p_T > 100 \text{ GeV}$  and subleading jets with  $p_T > 50 \text{ GeV}$
- 3.  $E_T > 100 \text{ GeV} \text{ and } E_T > 0.2 M_{\text{eff}}$
- 4. Transverse sphericity  $S_T > 0.2$ .
- Detector effects: Simple gaussian smearing

$$\frac{\Delta E_e}{E_e} = \frac{10\%}{\sqrt{E_e(\text{GeV})}} \oplus 0.7\%, \qquad \frac{\Delta E_j}{E_j} = \frac{50\%}{\sqrt{E_j(\text{GeV})}} \oplus 3\%.$$

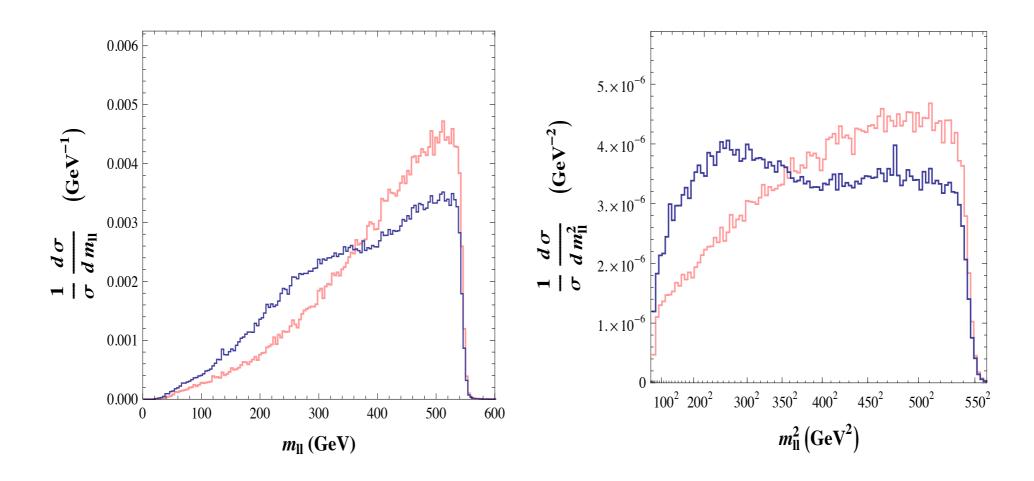
# Backgrounds

(in one transparency)

Background	Events $(1 \text{ fb}^{-1})$
$t\overline{t}$	81.5
W+ jets	1.97
Z+ jets	1.20
QCD	0
Total SM	84.67

 Additional signal background: Invisible Z decays

 Dilepton invariant mass (with additional diagrams and cuts)

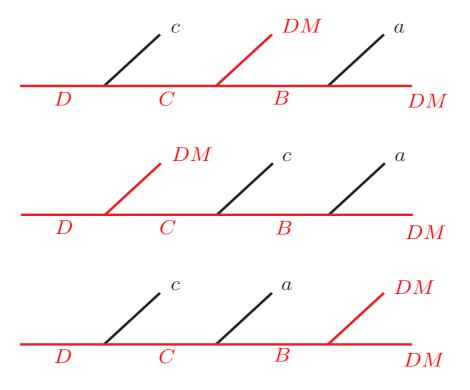


$$S/B = 0.08$$

$$S/\sqrt{B} = 18.6$$

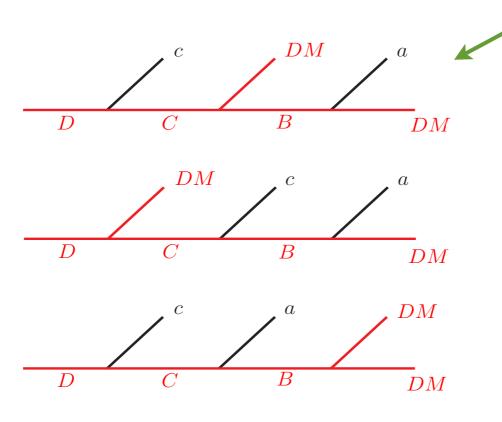
# Part II: Topologies and MT2 Reconstruction

So far focused on one topology



So far focused on one topology

Discussed this topology first.



• So far focused on one topology

Discussed this topology first.

Do C B DM

If these topologies exist, we want to account for them.

• So far focused on one topology

Discussed this topology first.

DM

DM

C

B

DM

If these topologies exist, we want to account for them.

 Emphasize: Models with "non-parity" stabilization symmetries can have more than one DM per leg. Use MT2 to tell the difference.

Basically the transverse mass:

$$\left(M_T^{(i)}\right)^2 = \left(m_T^{v(i)}\right)^2 + \tilde{m}^2 + 2\left(E_T^{v(i)}\tilde{E}_T^{(i)} - \mathbf{p}_T^{v(i)} \cdot \tilde{\mathbf{p}}_T^{(i)}\right)$$

Basically the transverse mass:

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Visible transverse mass

Basically the transverse mass:

Missing mass (or trial mass)  $\left(M_T^{(i)}\right)^2 = \left(m_T^{v(i)}\right)^2 + \tilde{m}^2 + 2\left(E_T^{v(i)}\tilde{E}_T^{(i)} - \mathbf{p}_T^{v(i)}\cdot\tilde{\mathbf{p}}_T^{(i)}\right)$ 

Visible transverse mass

Basically the transverse mass: Transverse energy

Missing mass (or trial mass)  $\left(M_T^{(i)}\right)^2 = \left(m_T^{v(i)}\right)^2 + \tilde{m}^2 + 2\left(E_T^{v(i)}\tilde{E}_T^{(i)} - \mathbf{p}_T^{v(i)} \cdot \tilde{\mathbf{p}}_T^{(i)}\right)$ 

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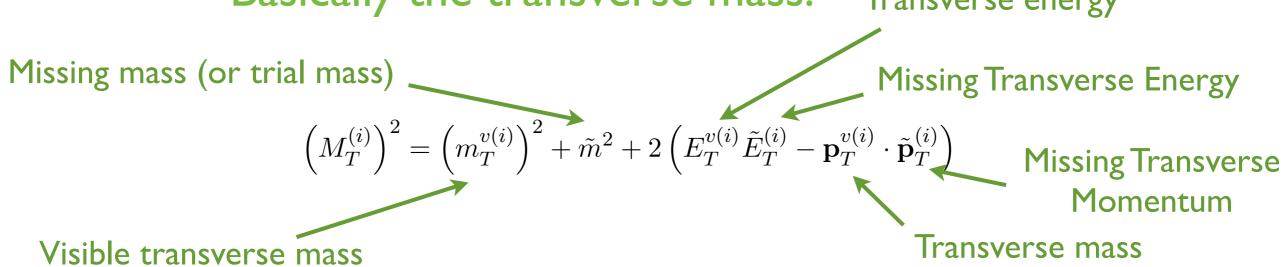
Visible transverse mass

Basically the transverse mass: Transverse energy

Missing mass (or trial mass)  $\left(M_T^{(i)}\right)^2 = \left(m_T^{v(i)}\right)^2 + \tilde{m}^2 + 2\left(E_T^{v(i)}\tilde{E}_T^{(i)} - \mathbf{p}_T^{v(i)}\cdot\tilde{\mathbf{p}}_T^{(i)}\right)$  Missing Transverse Momentum

Visible transverse mass

Basically the transverse mass: Transverse energy



Minimize transverse mass for each decay leg

$$M_{T2} \equiv \min_{\mathbf{p}_{T}^{v(1)} + \mathbf{p}_{T}^{v(2)} + \mathbf{p}_{T}^{v(2)} + \mathbf{p}_{T}^{v} = 0} \left[ \max \left\{ M_{T}^{(1)}, M_{T}^{(2)} \right\} \right]$$

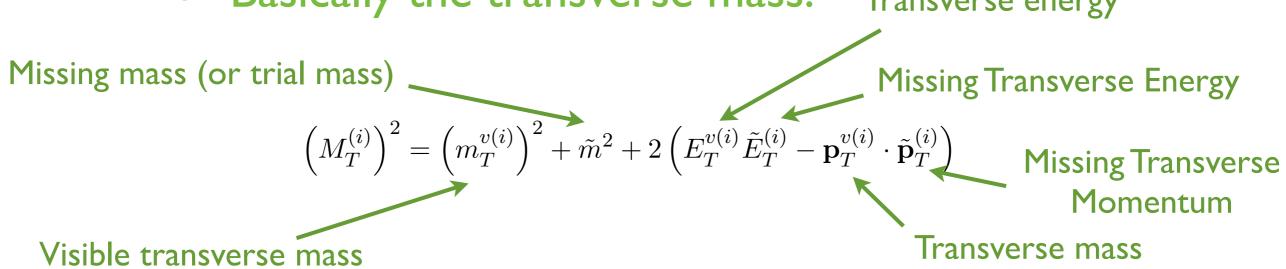
Minimize transverse mass for each decay leg

Transverse mass

Maximum value over 
$$many \ events \\ M_{T2} \equiv \min_{\substack{\mathbf{p}_T^{v(1)} + \mathbf{p}_T^{v(2)} + \mathbf{p}_T = 0}} \left[ \max \left\{ M_T^{(1)}, \ M_T^{(2)} \right\} \right] \\ M_{T2}^{\max}(\tilde{m}) = \max_{\substack{\text{many events}}} \left[ M_{T2}(\tilde{m}) \right]$$

Visible transverse mass

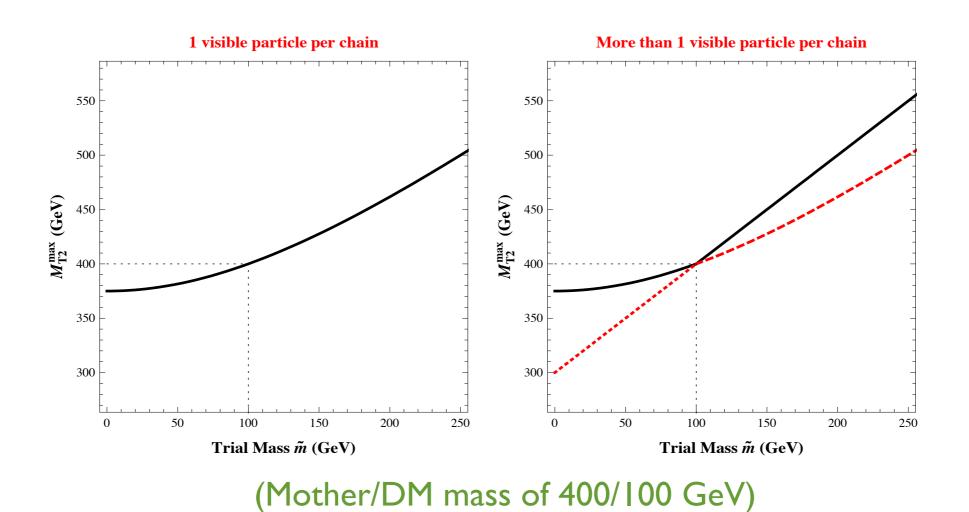
Basically the transverse mass: Transverse energy



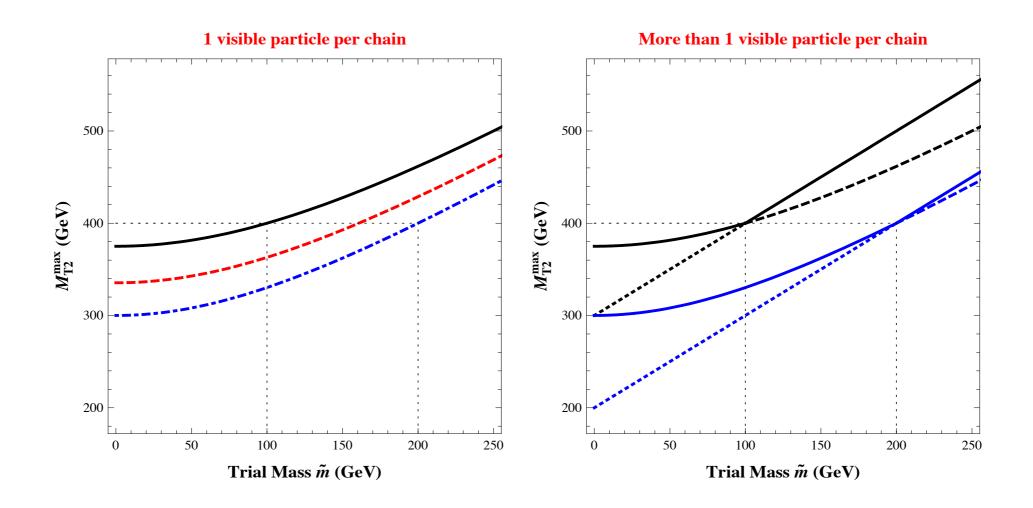
Minimize transverse mass for each decay leg

Maximum value over  $many \ \text{events}$   $M_{T2} \equiv \min_{\substack{\mathbf{p}_T^{v(1)} + \mathbf{p}_T^{v(2)} + \mathbf{p}_T = 0}} \left[ \max \left\{ M_T^{(1)}, \ M_T^{(2)} \right\} \right]$   $M_{T2}^{\max}(\tilde{m}) = \max_{\substack{\text{many events}}} \left[ M_{T2}(\tilde{m}) \right]$  Dependent on trial mass

• Generates a "kink" structure for Z2 models:

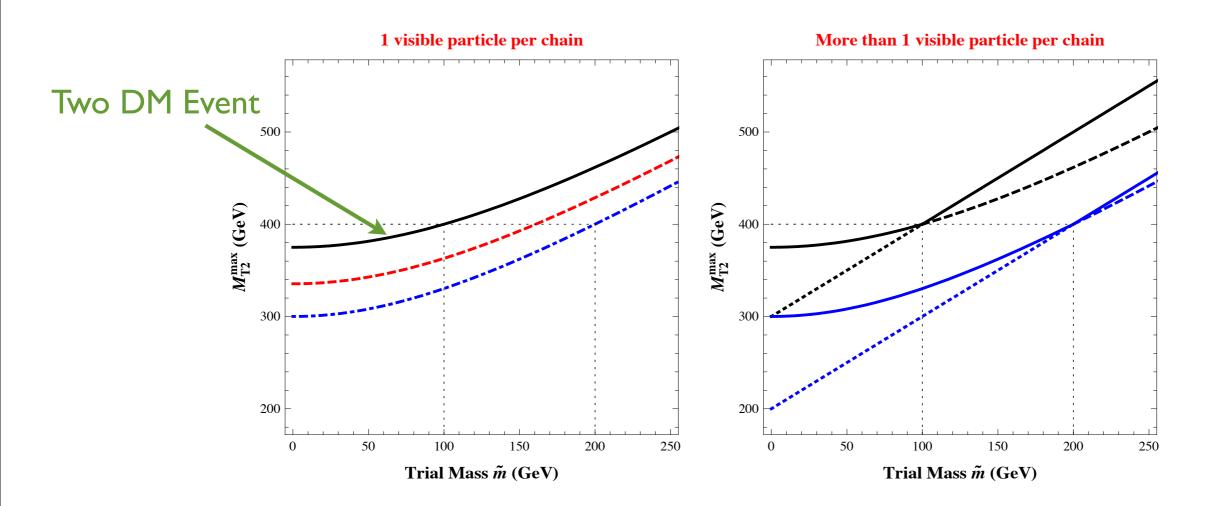


• Generates a "kink" structure for "non-parity" models!



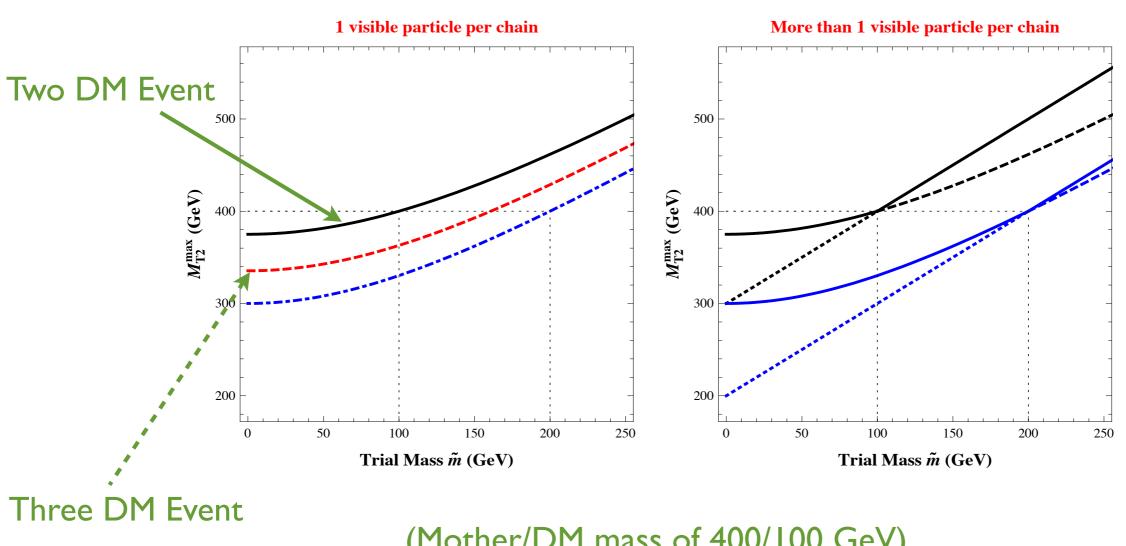
(Mother/DM mass of 400/100 GeV)
(Trial mass of 25 GeV)

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(Trial mass of 25 GeV)

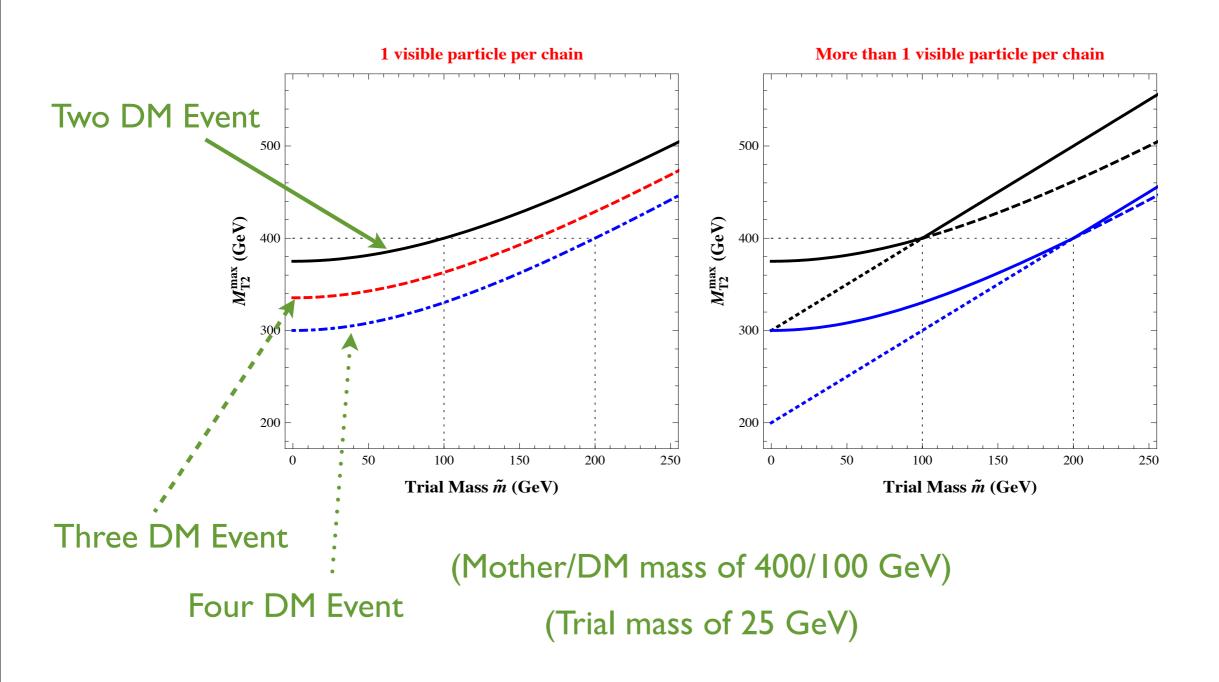
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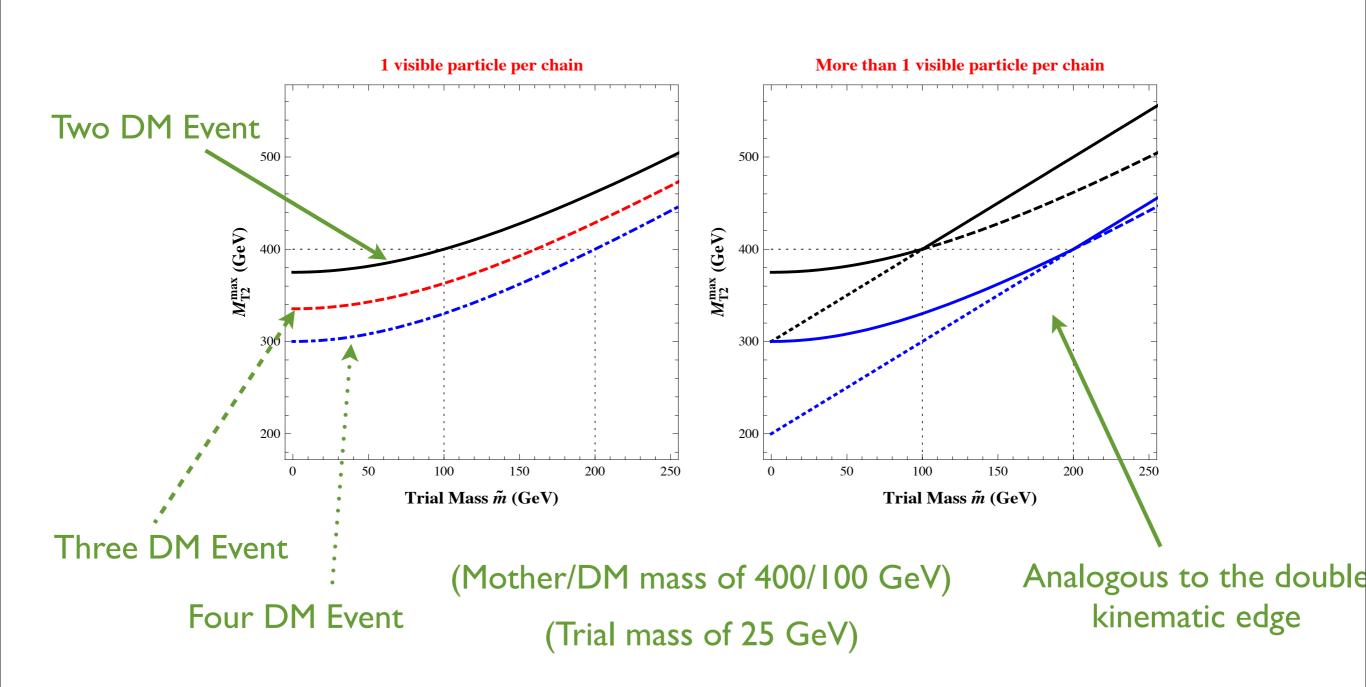
(Mother/DM mass of 400/100 GeV)

(Trial mass of 25 GeV)

• Generates a "kink" structure for "non-parity" models!



• Generates a "kink" structure for "non-parity" models!



 Highly dependent on properly reconstructing events. (Address this in the future).

## "Non-parity" and MT2

 Care required in seeing "non-parity signal" over Z2 bkg.

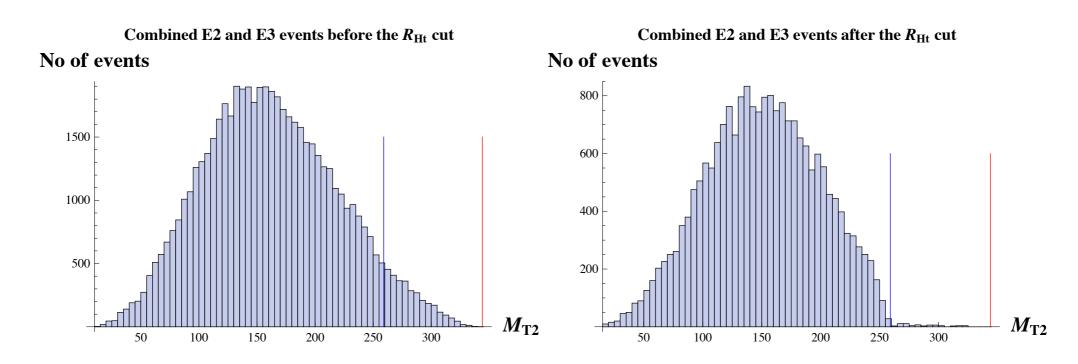
## "Non-parity" and MT2

- Care required in seeing "non-parity signal" over Z2 bkg.
  - Define a ratio of the sum of the visible pT on each leg.

$$H_t^i = \sum_a |P_t^{v_a^i}|$$

$$R_{H_t} = \frac{H_t^{\text{max}}}{H_t^{\text{min}}} \qquad H_t^{\text{max}} = \max(H_t^1, H_t^2) \\ H_t^{\text{min}} = \min(H_t^1, H_t^2)$$

# "Non-parity" Signal

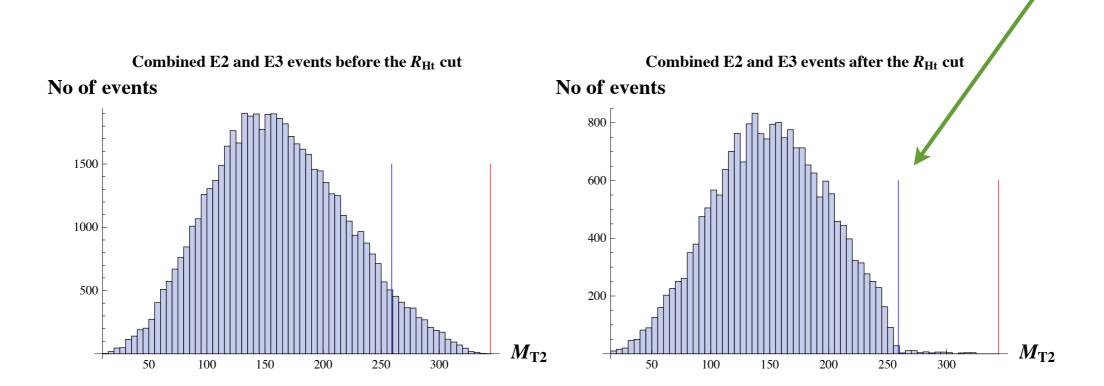


(Mother/DM mass of 400/150 GeV)

(Ht > 5/Trial mass of 9 GeV)

# "Non-parity" Signal

MT2 predicted upper edge for three DM in event



(Mother/DM mass of 400/150 GeV)

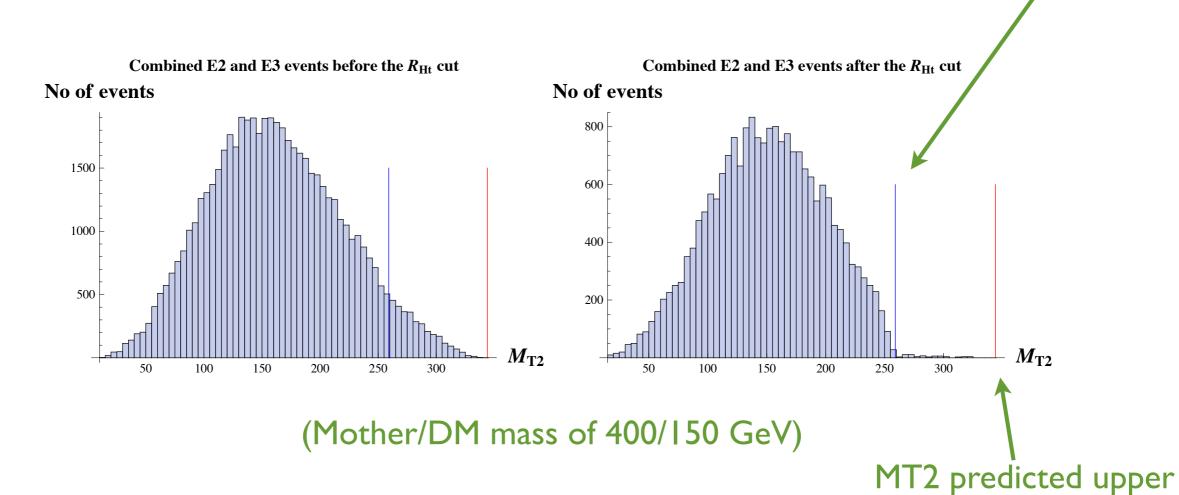
(Ht > 5/Trial mass of 9 GeV)

# "Non-parity" Signal

MT2 predicted upper edge for three DM in event

edge for two DM

in event



(Ht > 5/Trial mass of 9 GeV)

# Part III: Signatures with Metastable Particles

#### What is a Meta-Stable Particle?

- Do not decay in the LHC's detectors.
- Charged under the SM and dark matter stabilization symmetry.
- Lifetime consistent with known bounds.\*

\* Perl, Kim Halyo, Lee, Lee, Loomba and Lackner, Int. J. Mod. Phys. A16 2137 (2001).

## Key Points

• Signal: Final states with two meta-stable particles plus large missing energy at the LHC.

## **Key Points**

• Signal: Final states with two meta-stable particles plus large missing energy at the LHC.

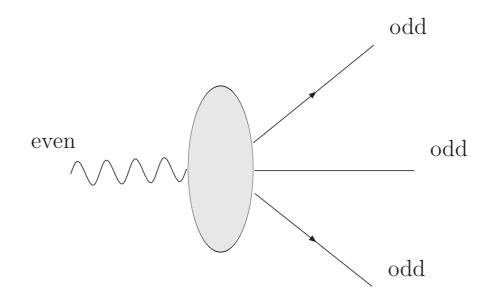
 Claim: Models with parity stabilized dark matter are generally suppressed and distinguishable from models with "nonparity" dark matter.

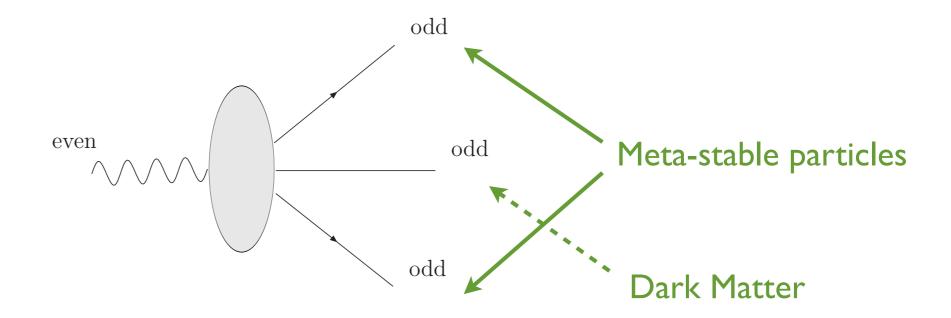
## **Key Points**

• Signal: Final states with two meta-stable particles plus large missing energy at the LHC.

 Claim: Models with parity stabilized dark matter are generally suppressed and distinguishable from models with "nonparity" dark matter.

 Strategy: Focus on parity stabilized models to see the differences.





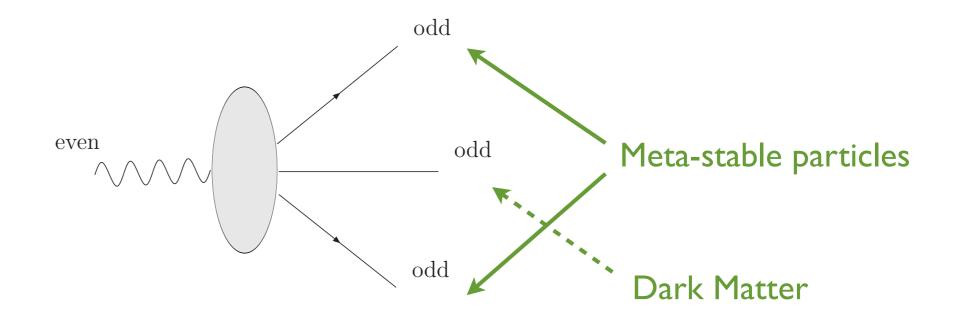
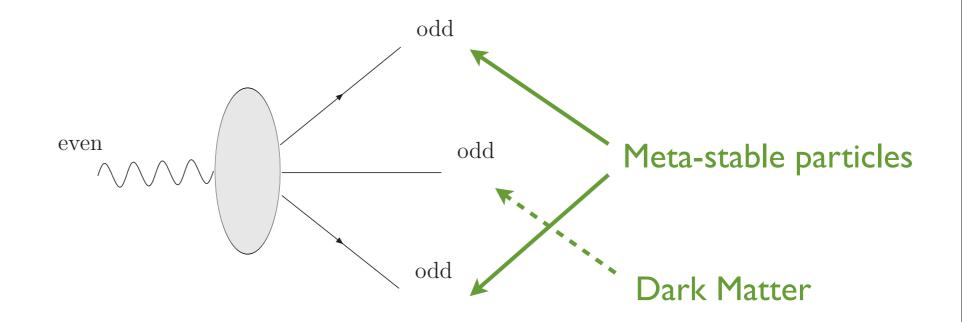
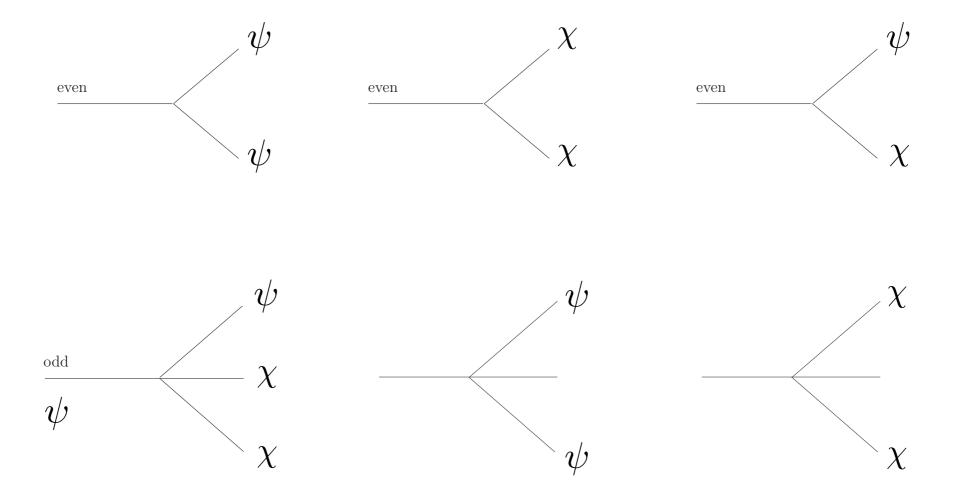


 Diagram rarely happens even for final states with two metastable particles plus an even number of DM.

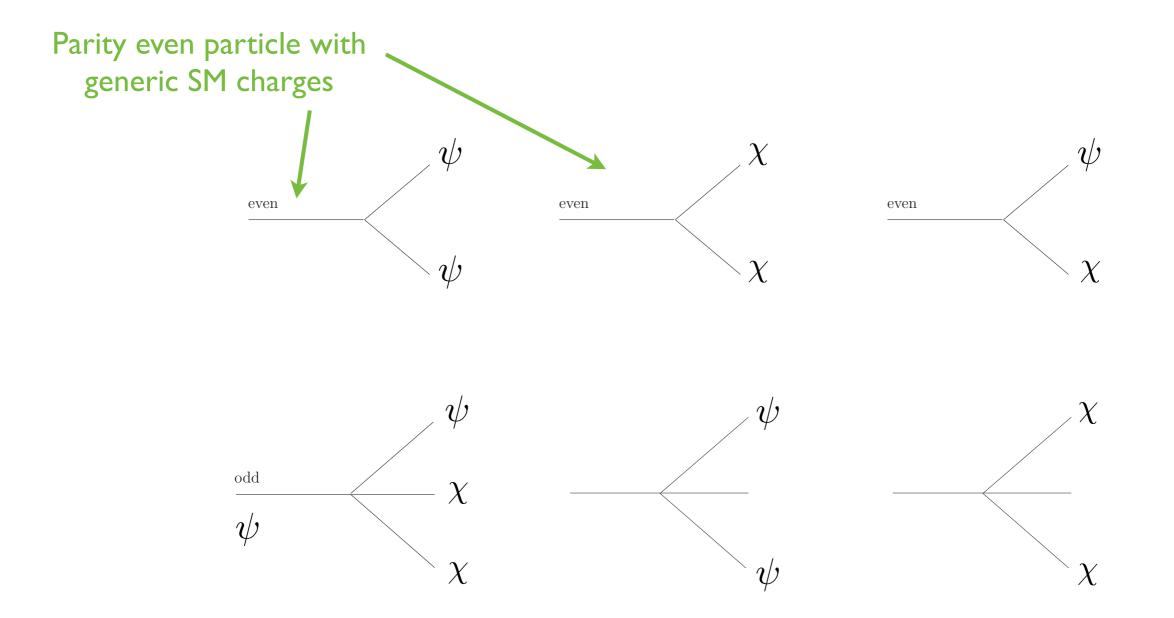


- Diagram rarely happens even for final states with two metastable particles plus an even number of DM.
- Add new particles to prove this...

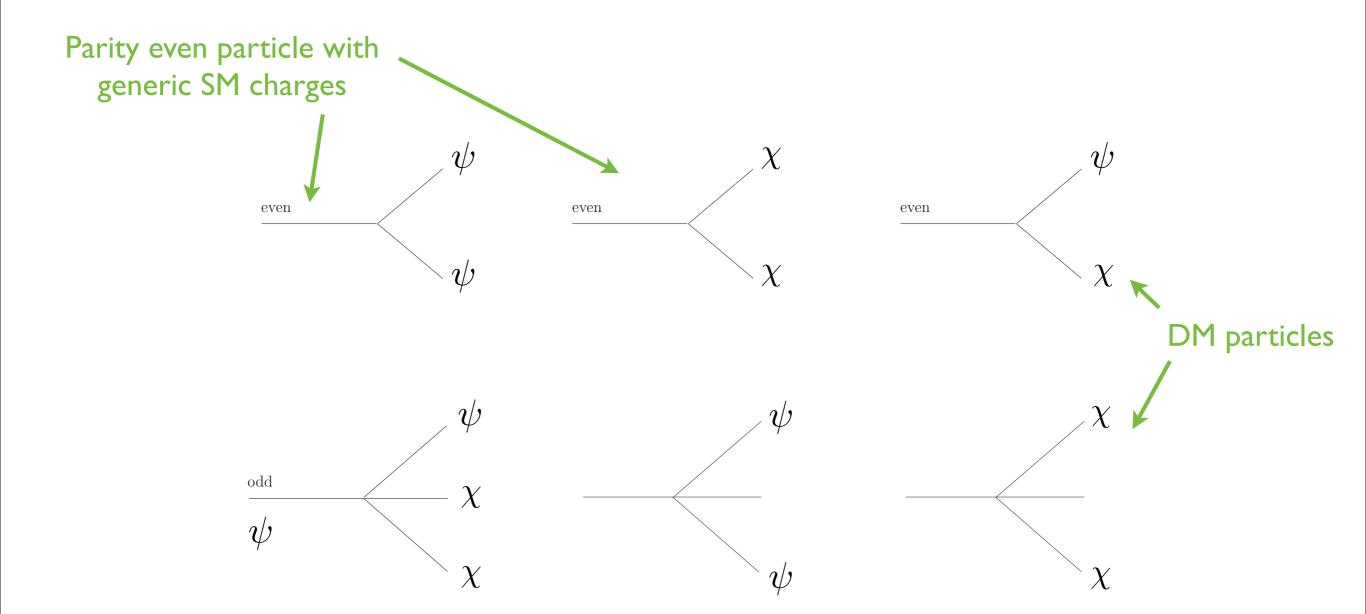
## All Possible Relevant Couplings



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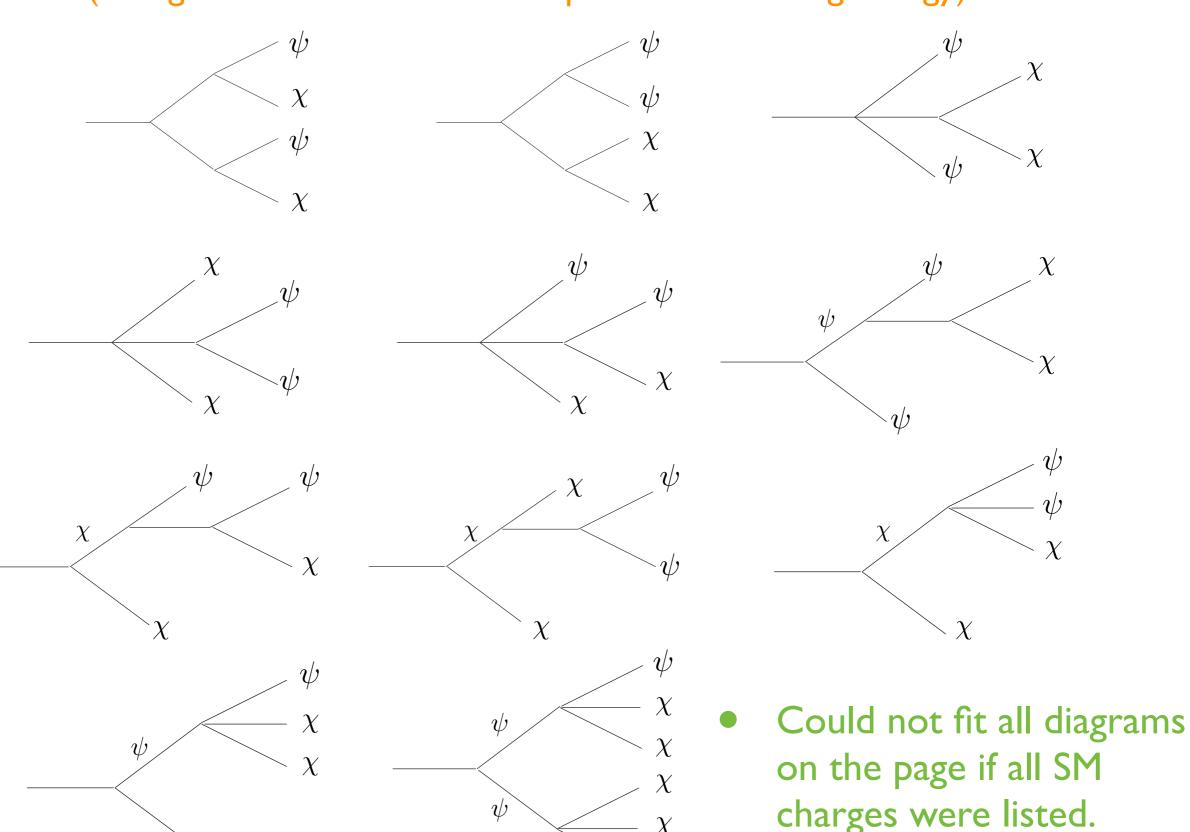


## All Possible Relevant Couplings



### All Possible Four Body States

(that generate two meta-stable particles + missing energy)



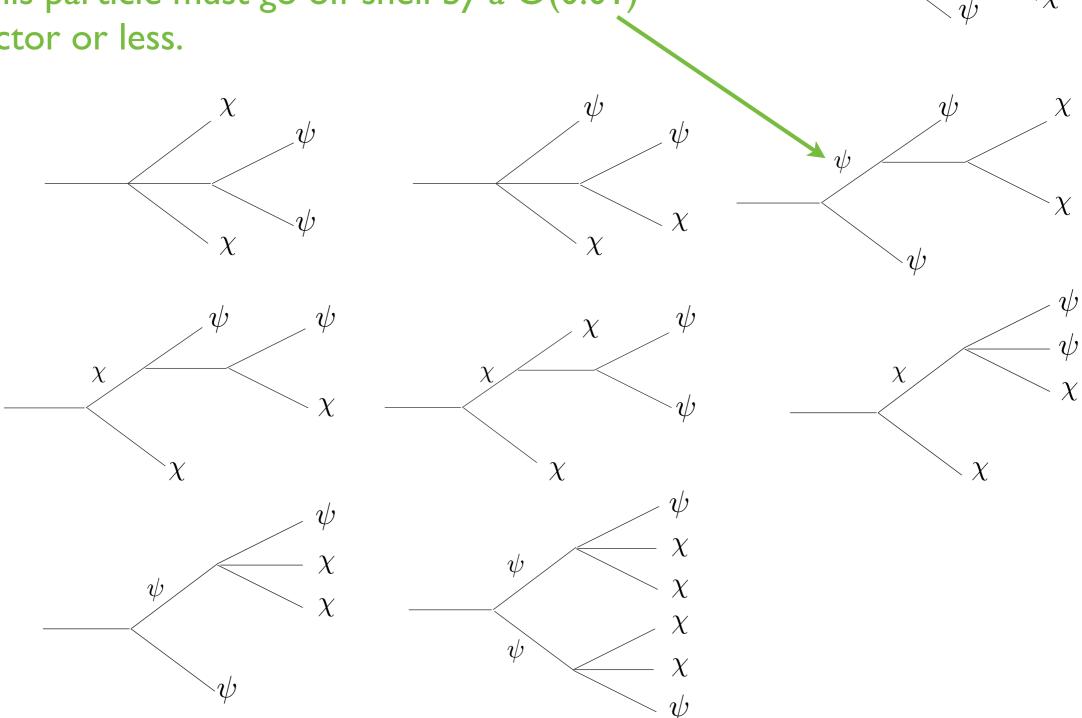
Thursday, October 7, 2010

## Off Shell Suppression

 $\chi$ 

All of these diagrams have off-shell suppression...

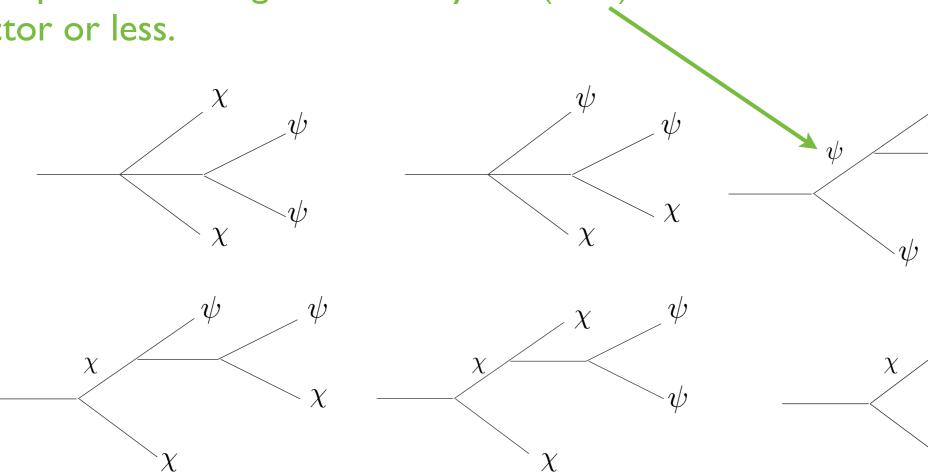
Consider 500 GeV  $\psi$  and a 100 GeV  $\chi$ . This particle must go off-shell by a O(0.01) factor or less.

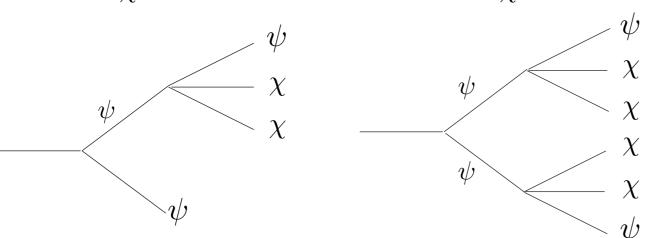


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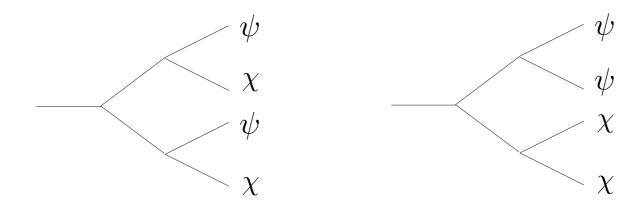
Strong SuppressionConstraint

 $\chi$ 

 $\psi$ 

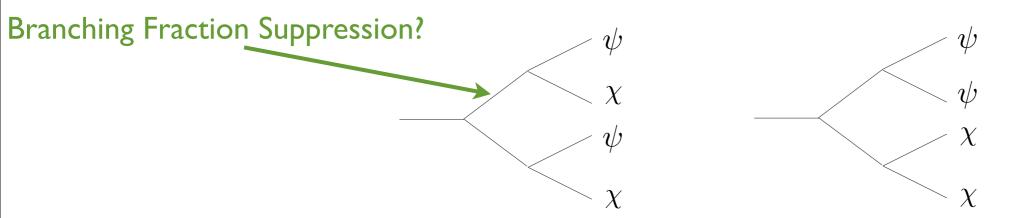
 $\chi$ 

# Remaining Diagrams



• Constraints on the remaining diagrams?

## Remaining Diagrams?



An example: Virtual SM higgs particle. Decay rate

$$\Gamma_{h_0 \to \chi \chi \text{ scalar}} = \frac{\kappa_1^2 v_{ew}^2}{16\pi m_{h_0}} \sqrt{1 - 4m_{\chi}^2 / m_{h_0}^2}$$

$$\Gamma_{h_0 \to \chi \chi \text{ fermion}} = \frac{A \kappa_2^2 m_{h_0}}{32\pi} \left(1 - 4m_{\chi}^2 / m_{h_0}^2\right)^{3/2}$$

I TeV SM higgs branching fraction ~ O(0.01)% 350 GeV higgs branching ~ O(1)% (largest possible)

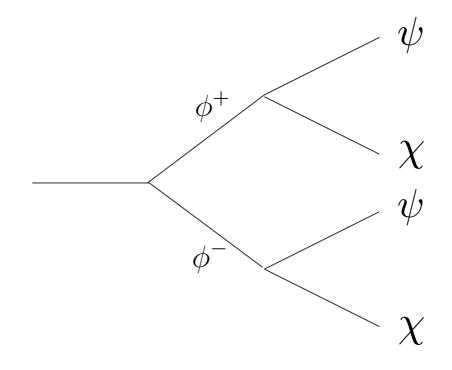
## Remaining Diagrams Suppressed?

• Arguments are not 100%.

## Remaining Diagrams Suppressed?

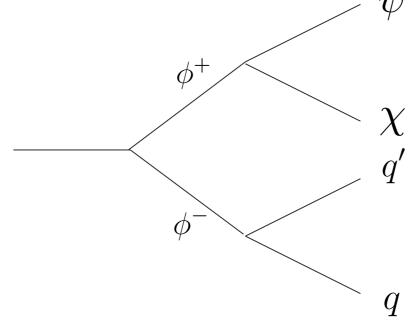
• Arguments are not 100%.

 Best we can do? Try to reconstruct the new scalar to see if it is parity even.



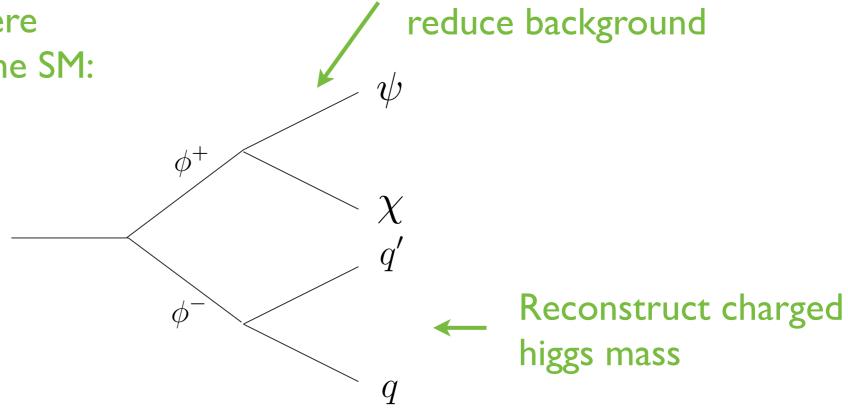
Search for processes where one branch decays into the SM:

$$p p \rightarrow \psi \chi^* + SM$$



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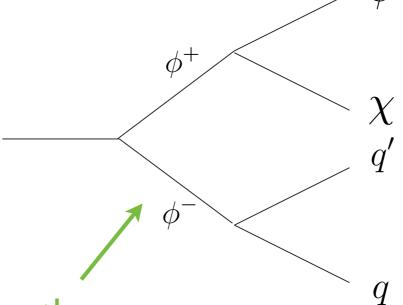
$$p p \rightarrow \psi \chi^* + SM$$



Tag meta-stable to

Search for processes where one branch decays into the SM:

 $p p \rightarrow \psi \chi^* + SM$ 



Reconstructing the mass gives us a "tag" and verifies this is parity even.

Tag meta-stable to reduce background

Reconstruct charged higgs mass

Search for processes where one branch decays into the SM:

 $p p \rightarrow \psi \chi^* + SM$ 

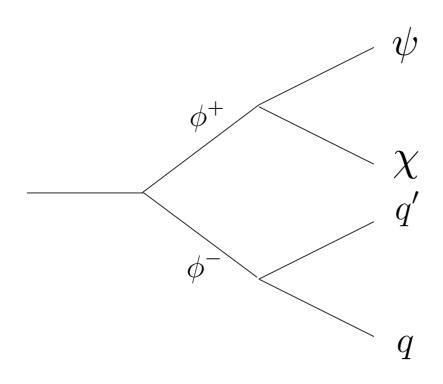
 $\begin{array}{c} \phi^+ \\ \chi \\ q' \\ \end{array}$ 

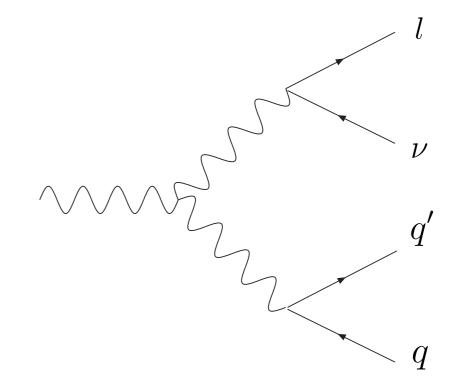
Tag meta-stable to reduce background

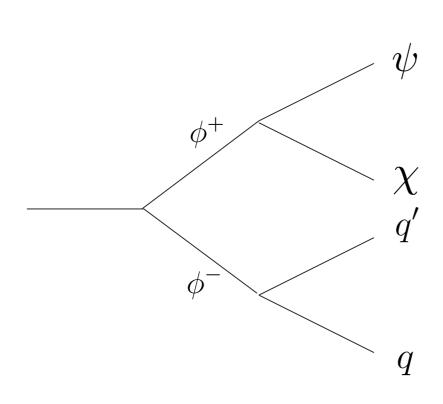
Reconstruct charged higgs mass

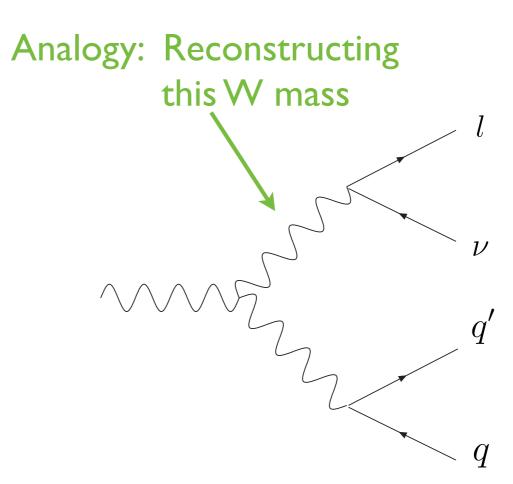
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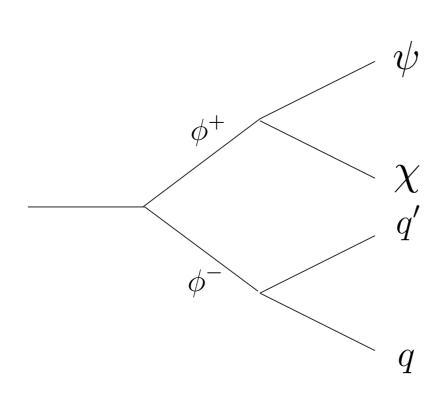
• Reconstruct in analogy to WW reconstruction to get the mass of  $\phi^+$ .

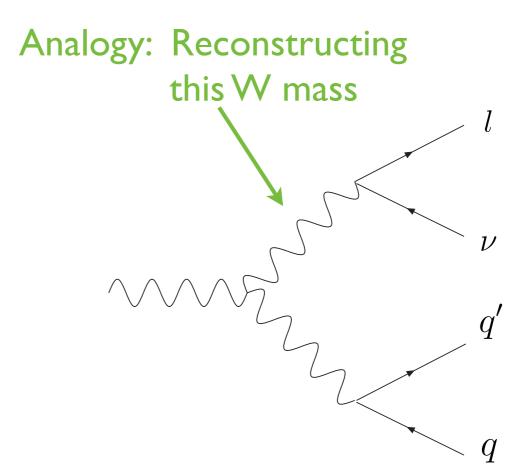








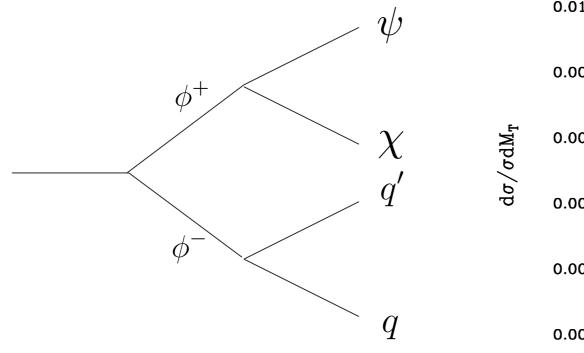


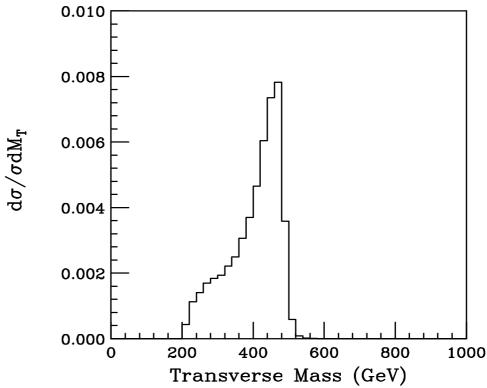


• Define variables: 
$$M_T = \sqrt{(E_T + E_T)^2 - (\vec{p}_T + \vec{p}_T)^2}$$

$$E_T^2 = \vec{p}_T^2 + M_\psi^2$$

$$E_T^2 = \vec{p}_T^2 - \vec{p}_T = \sum_{\tau} \vec{p}_{\tau} \text{ visible}$$

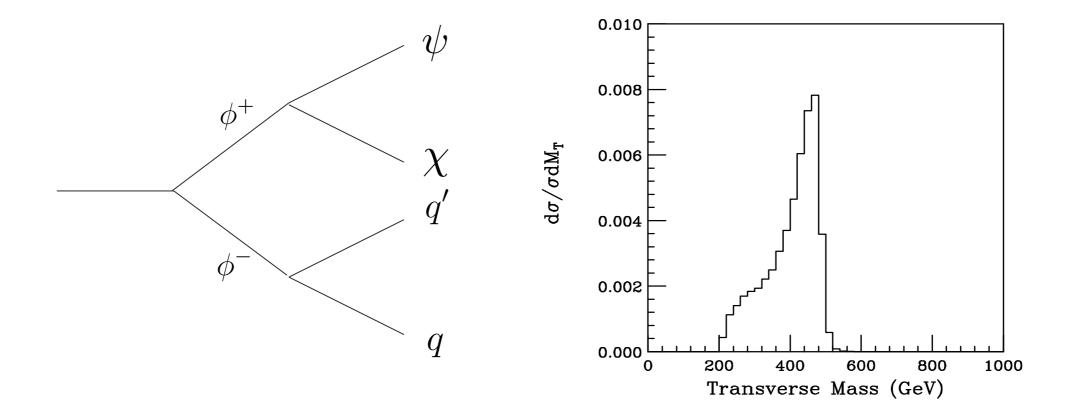




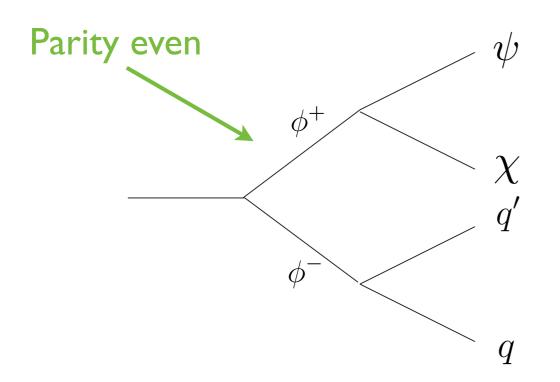
• Define variables: 
$$M_T = \sqrt{(E_T + \not\!\!E_T)^2 - (\vec{p}_T + \not\!\!P_T)^2}$$

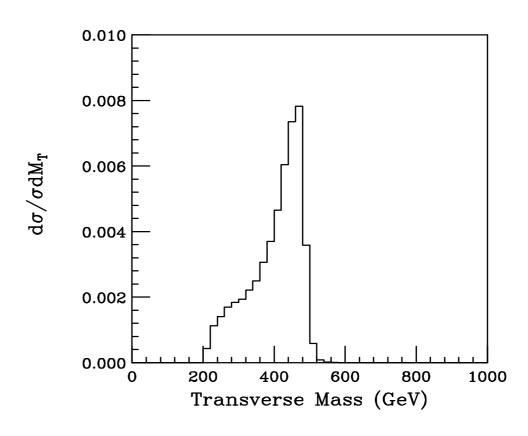
$$E_T^2 = \vec{p}_T^2 + M_\psi^2$$

$$\not\!\!E_T^2 = \not\!\!P_T^2 - \not\!\!P_T = \sum \vec{p}_{T \text{ visible}}$$



• 500 GeV  $\phi$ , 200 GeV  $\psi$  and 100 GeV DM  $\chi$ .

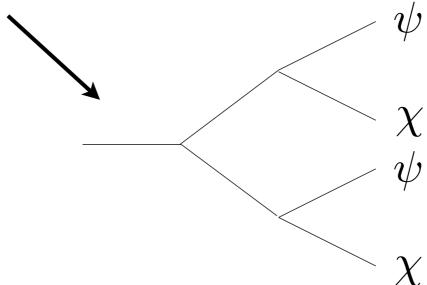




- 500 GeV  $\phi$ , 200 GeV  $\psi$  and 100 GeV DM  $\chi$ .
- Conclusion: Parity even scalar is generating signal.

## Going Forward

 We assume no scalars in our EFT which mediate processes like



• Onto collider studies ...

#### Model Parameters

• Take masses of the meta-stable particles to be

$$m_{\psi} = 300, 600 \text{ GeV}$$

All dark matter candidates are

$$m_{\chi} = 100 \text{ GeV}$$

### Backgrounds

- Many SM process produce a potential background: light quark production (QCD), bbbar, ttbar, W, Z, WW and ZZ.
- Any SUSY process with final state neutrinos.
- ATLAS and CMS have unique approaches to eliminate the backgrounds.

## An Additional SM Background

 SM can produce a similar missing energy process by radiating off a Z boson:

$$p\, p \to Q\, \overline{Q}\, Z \to Q\, \overline{Q}\, \nu\, \overline{\nu} \qquad \qquad q\, \overline{q} \to L\, \overline{L}\, Z \to L\, \overline{L}\, \nu\, \overline{\nu}$$

Important the signal cross section is no less than 10% of the SM background (conservatively) to preclude statistical fluctuations

### An Additional SM Background

- Some numbers: (with 300 GeV "hadron" no cuts)
  - Cross section to radiate off a Z from heavy stable quarks and decay invisibly:  $0.33 \text{ pb} \times 0.2 = 0.07 \text{ pb}$ .
  - Cross section to radiate off a 100 GeV heavy gauge boson: 0.28 pb.
  - Cross section to radiate off a 275 GeV heavy gauge boson: 0.01 pb.
- No problem with this background.

#### ATLAS Cuts

- ATLAS requires the "muon tracks" to have a transverse momentum of  $p_T > 135~{\rm GeV}$  for a 300 GeV "hadron."
- The events are restricted to the "triggerable" part of the muon detector  $|\eta| \le 2.4$ .
- No jet in the inner detector must come within a cone of  $\Delta R = < 0.4$  of the "muon track."
- The "muon tracks" must have a velocity between  $0.7 < \beta_{\rm ATLAS} < 0.9$  .

\*See, for example, hep-ex/05011014

#### **CMS** Cuts

 CMS detector muon detector is partly suspended by iron yoke.

The "hadrons" can charge flip even in the muon detector.

- The collaboration requires a highly ionizing track in the inner detector.
- The analysis: Kept events with charged "hadrons" in the inner detector.

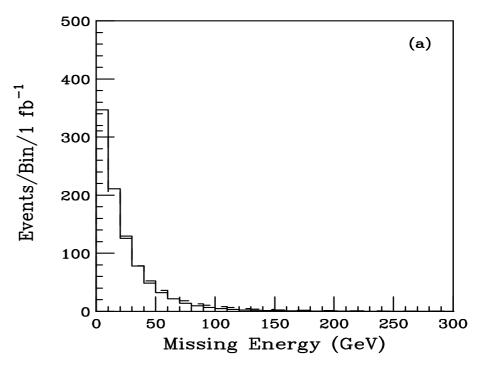
### CMS Cuts\*

- CMS require an invariant mass of  $m>100~{\rm GeV}$ .
- CMS requires the "muon tracks" to have a velocity of  $0.45 < \beta_{\rm CMS} < 0.8$ .
- For these cuts, CMS finds "a background free region."
- Since the S/B is so high, we safely neglect the background.

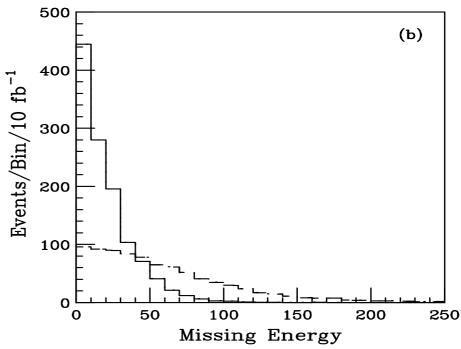
\*CMS Collaboration, CMS PAS EXO-08-003

# Signal from Acceptance Cuts?

Atlas (solid) and CMS (dotted) for the signal production (chalkboard).



Same plot as above but with ATLAS and CMS cuts.



• To clarify signal we require an additional cut.

### "Non-parity" Results

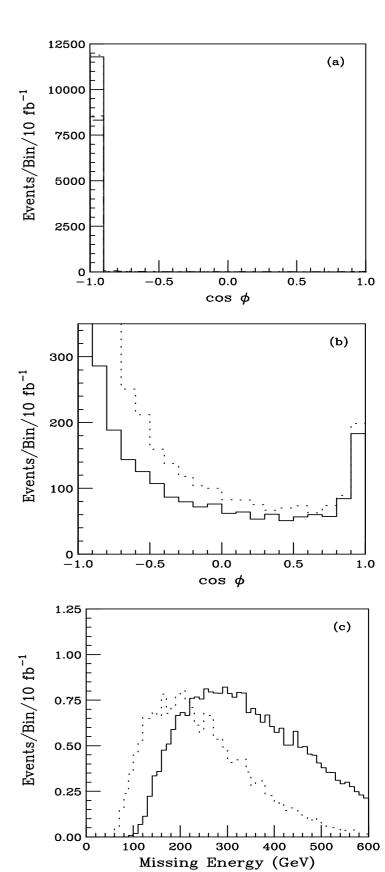
Transverse angle for ATLAS (solid) and CMS(dotted). Most events are back-to-back. All plots use feature 100 GeV DM and 300 GeV meta-stable particles. All kinematic cuts are used.

Same plot as above but emphasizing the non back-to-back events.

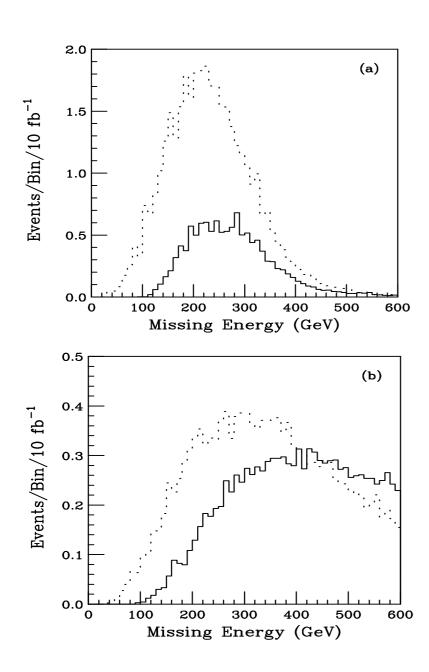
Applying the cut

$$\cos \phi > -0.9$$
.

None of the back-to-back events survived. About 70 events of signal for 10 inverse fb.



### More Results



Same plots with all cuts for the Agashe-Servant scenario (a) and light hidden sector (b).

#### Companion Model

-	ATLAS		CMS	
	$S/\sqrt{B}$	S/B	$S/\sqrt{B}$	S/B
$300~{\rm GeV}~\psi$	34.3	4.1	29.4	3.92
$600~{ m GeV}~\psi$	10.1	4.1	9.5	4.1

#### Agashe-Servant

	ATL	AS	CMS		
	$S/\sqrt{B}$	S/B	$S/\sqrt{B}$	S/B	
$300~{\rm GeV}~\psi$	34.3	4.1	29.4	3.92	

#### Light Hidden Sectors

	ATLAS		CMS		
	$S/\sqrt{B}$	S/B	$S/\sqrt{B}$	S/B	
$300 \text{ GeV } \psi$	5.33	0.63	8.76	1.04	
$600~{\rm GeV}~\psi$	0.75	0.30	1.08	0.47	

### Conclusion/Future Goals

If the LHC is relevant for dark matter hard questions like "what is the stabilization symmetry of dark matter" may be possible.

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If the LHC is relevant for dark matter hard questions like "what is the stabilization symmetry of dark matter" may be possible.

Implement ideas on reconstructing multi-hard jet events with timing.

Look to distinguish strongly coupled hidden sectors from perturbative ones...



#### **A Simulation**

- Need to reliably estimate...
  - How many "hadrons" charge flip to neutral? (Neutral "hadrons" generate obscuring missing energy.)
  - How much missing energy comes with and without emission of massive dark gauge bosons?
  - What are the cross sections?

#### **A Simulation**

- More things to reliably estimate...
  - How many of the stable particles stop in the detector generating more obscuring missing energy?

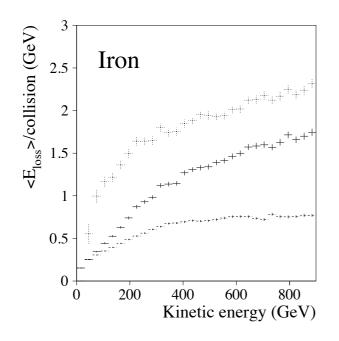
Energy fraction of the exchanged SM quark partons is  $\mathcal{O}(0.001)$  of the "hadron" total.\* Most interactions are low-energy QCD. Very few "hadrons" are stopped; but need to account for the relatively small signal.

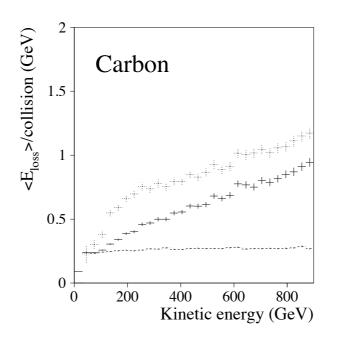
 GEANT4 without introduction is very difficult to use. No access to collaboration software.

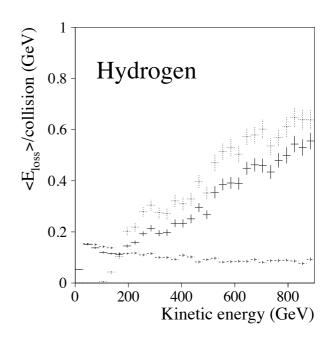
\*See review, Fairbairn, et al. hep-ph/0611040

### A Simulation

- Wrote a fast simulation of a perfect lead calorimeter that parameterizes the GEANT3 and 4 response to the interactions of "hadrons" with the calorimeters.\* Nuclear interactions only.
- Example response for 300 GeV R-hadron:



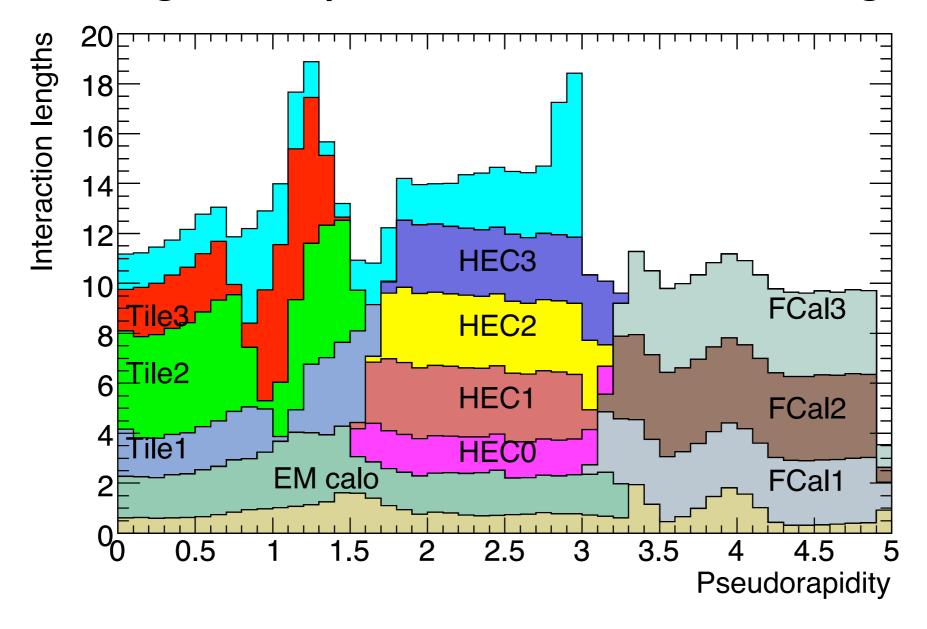




\*Kraan, hep-ph/0404001

### Detector Geometry

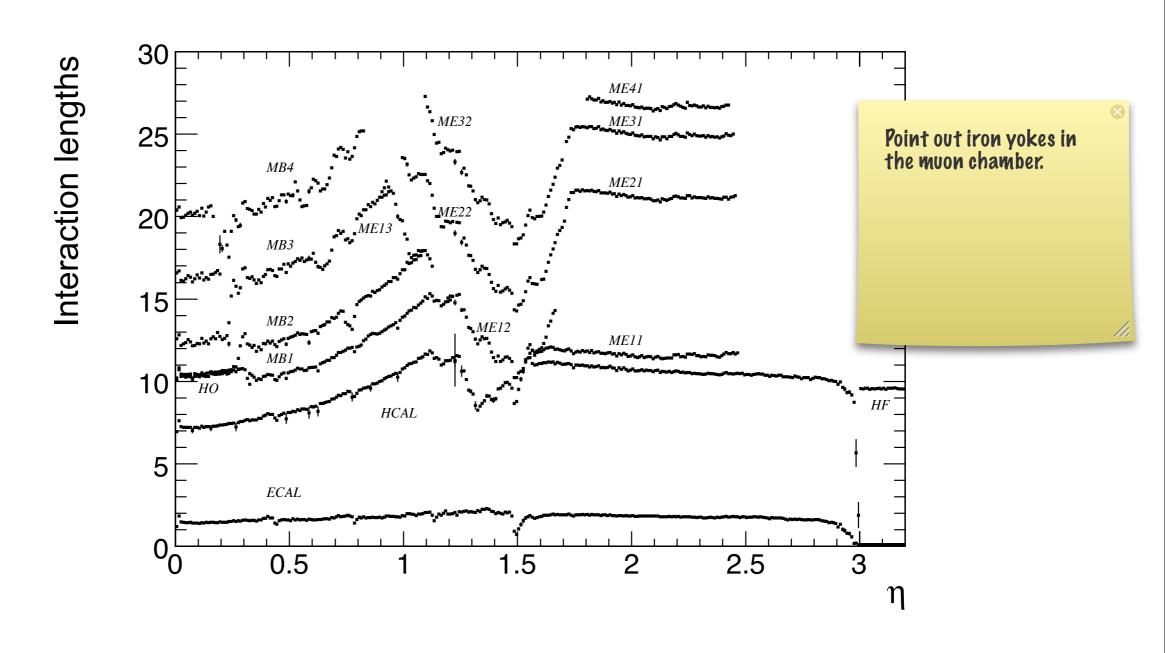
ATLAS geometry in nuclear interaction lengths\*



\*See G.Aad, et al. JINST 3, S08003 (2008)

### CMS Detector Geometry

CMS geometry in nuclear interaction lengths\*\*



\*See CMSTDR, Volume I

### Model Types

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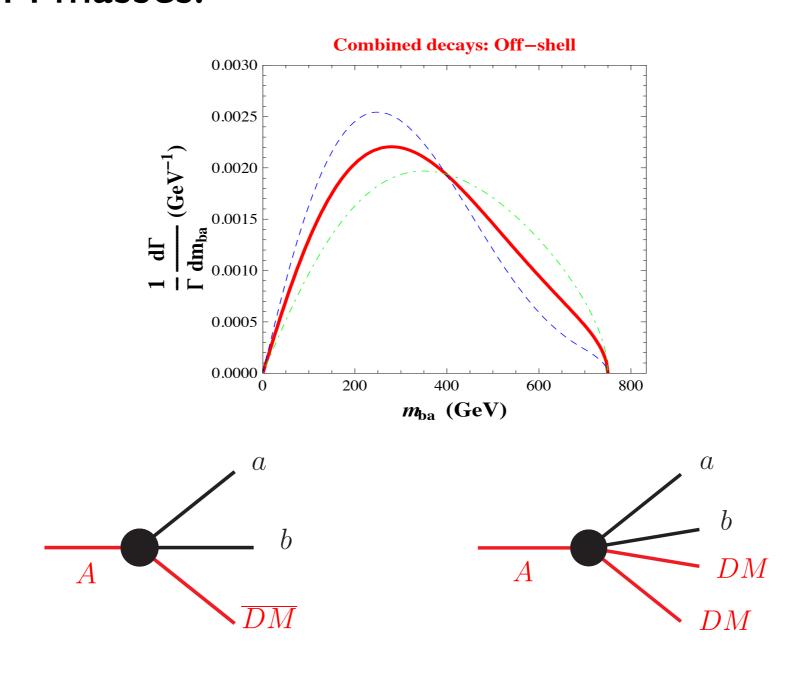
- Type II: DM symmetry associated with spontaneous symmetry breaking.
  - Focus: Model that generates the DM stabilization symmetry by SSB thereby correlating the DM relic abundance to the weak scale.\*
  - Also discuss how modifications to models which have non-trivial dark sectors can generate a signal.\*\*

\*D.W., arXiv:0907. 3146.

<sup>\*\*</sup>Arkani-Hamed, Finkbeiner, Slatyer and Weiner, PRD 79, 015014 (2009), Arkani-Hamed and Weiner, JHEP 0812, 104 (2008), Pospelov, Ritz and Voloshin, PLB 662, 53 (2008).

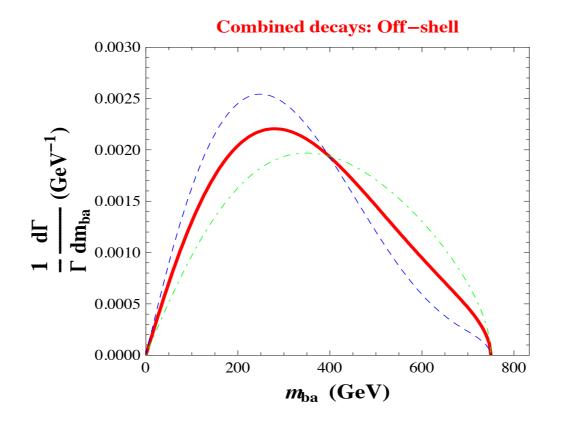
### Test

 Effect for large mass differences and small DM masses:



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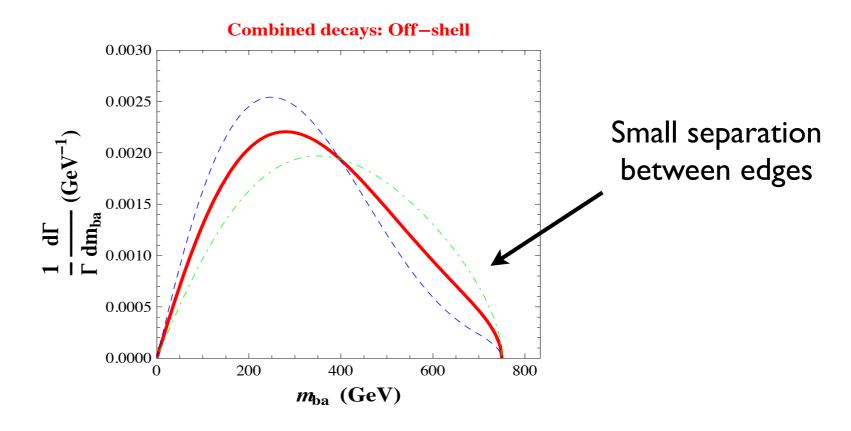
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 800 GeV mother and 50 GeV DM with no spin correlations.

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 800 GeV mother and 50 GeV DM with no spin correlations.

## **Model Types**

• Type I: Warped Pati-Salam Unification models where the DM symmetry is a  $\mathcal{Z}_3$ .\*

\*Agashe and Servant, PRL 93, 231805 (2004); JCAP 0502, 002 (2005).

## **Model Types**

• Type I: Warped Pati-Salam Unification models where the DM symmetry is a  $\mathcal{Z}_3$ .\*

•  $\mathcal{Z}_3$  symmetry is generated from color and fractional baryon number of each particle.

$$\eta \to \exp\left[2\pi i\left(B - \frac{n_c - \bar{n}_c}{3}\right)\right]\eta.$$

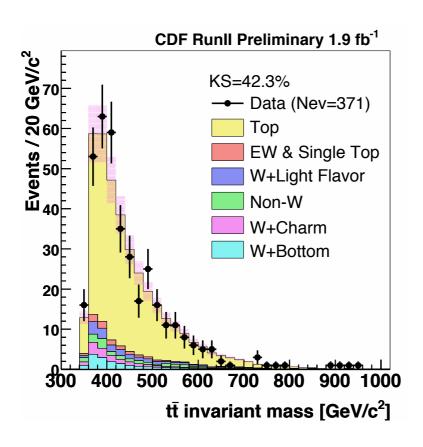
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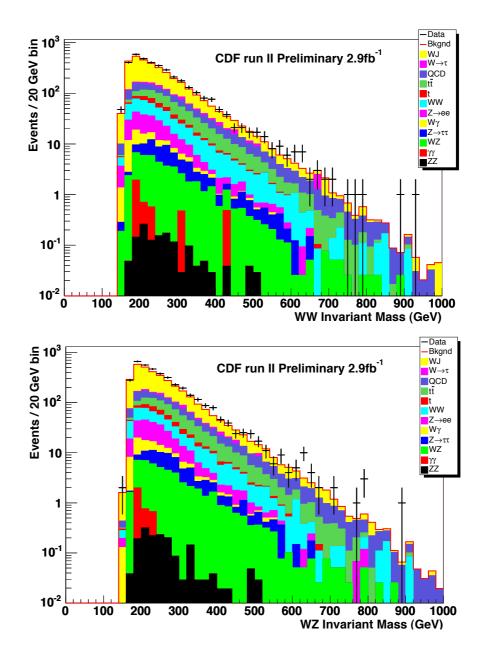


# Are the Diagrams Suppressed?

 TeVatron direct searches for parity even resonances exclude to about ~ 600 GeV.

**Examples:** 

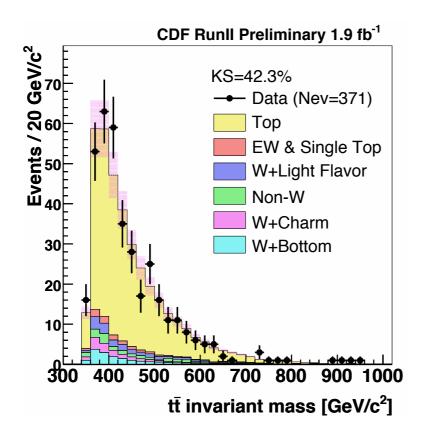


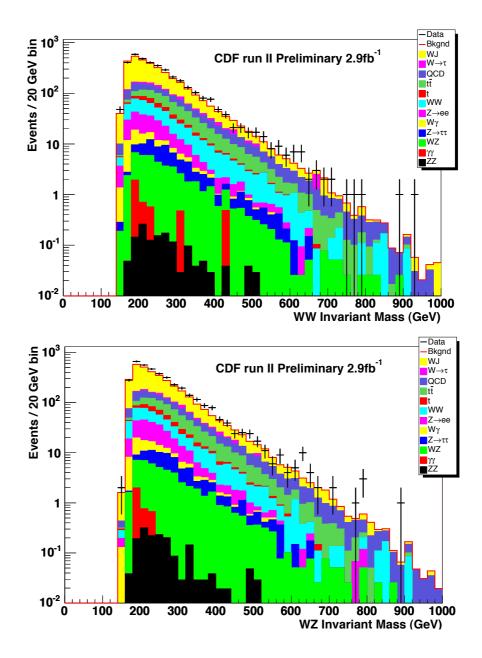


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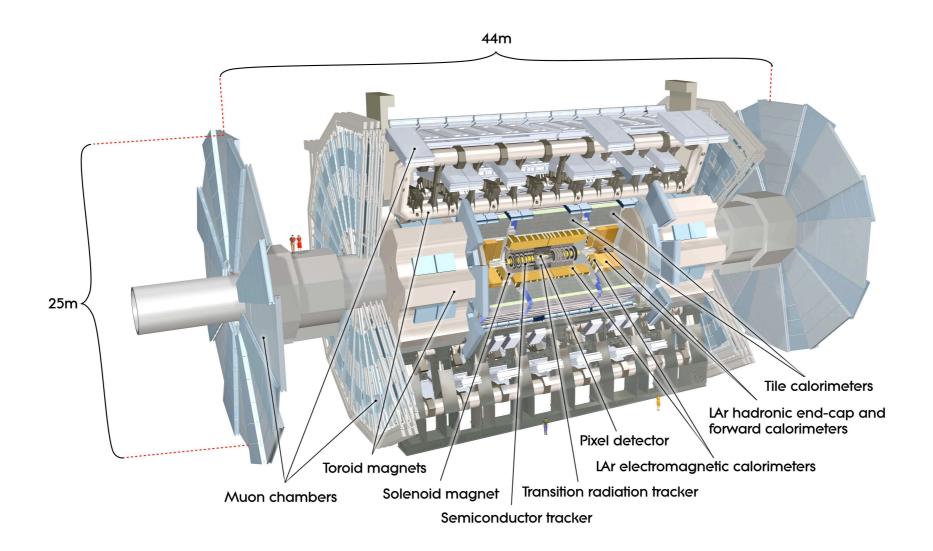
**Examples:** 





Still not strong.

## ATLAS Detector



## Arguments Void?

 LHC discovers scalars that have a large enough branching fraction into the signal such that a SM decay is very rare and some symmetry is not responsible.

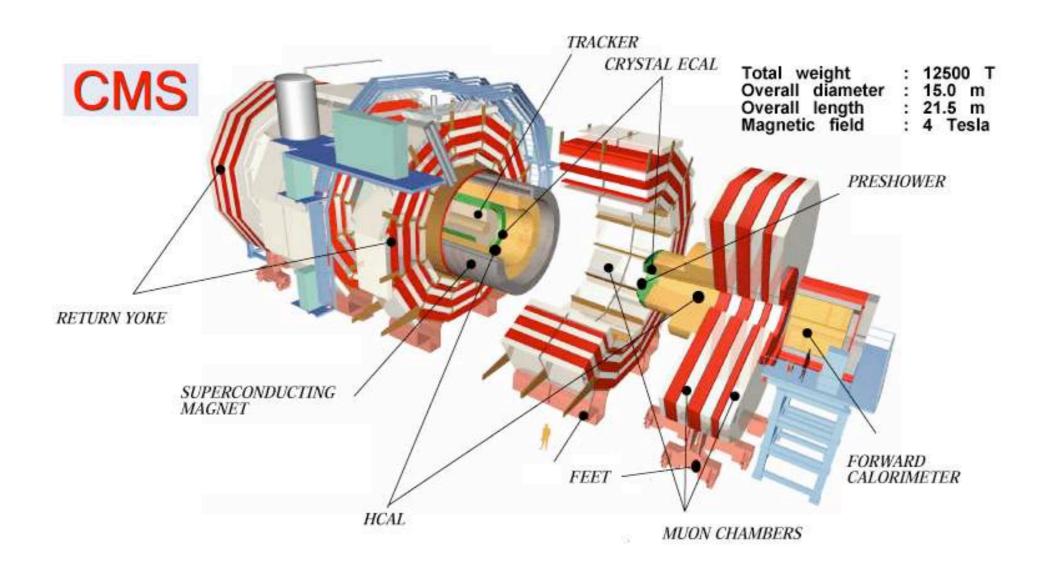
### Arguments Void?

 LHC discovers scalars that have a large enough branching fraction into the signal such that a SM decay is very rare and some symmetry is not responsible.

 Anthropic arguments? Fine-tune couplings of scalar w/SM to generate 100% fraction to signal.

Tantamount to adding a new symmetry to the DM. "Non-parity!"

# CMS Detector





## Assumptions

- Effective Field Theory: Start with the SM, DM  $(\chi)$  and meta-stable particle  $(\psi)$  below a cutoff.
- Long lived particles must decay via, e.g.,

$$\frac{1}{\Lambda^2} q q \psi \chi \qquad \frac{1}{\Lambda^2} l \nu \psi \chi$$

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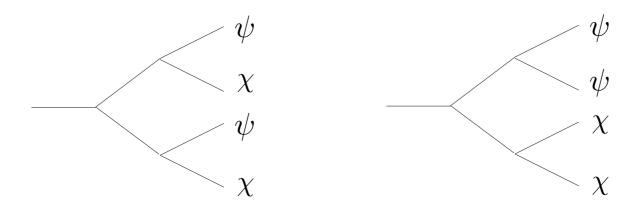
$$\frac{1}{\Lambda^2} q q \psi \chi \qquad \frac{1}{\Lambda^2} l \nu \psi \chi$$

Consequences:

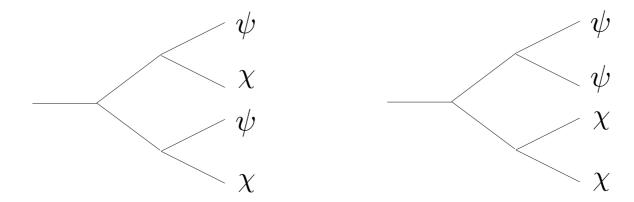
DM and meta-stable stabilization symmetries are correlated.

Cannot randomly pick non-SM charges for  $\psi$ .

# Remaining Diagrams Suppressed?



## Remaining Diagrams Suppressed?



 Effective Field theory: New particles generally have O(I) couplings to SM.

(remember new particles are parity even)

• Result: Smaller branching fraction into signal.

# Example Particle to Mediate Diagrams

An Estimate: Virtual SM higgs particle. Decay rate

$$\Gamma_{h_0 \to \chi \chi \text{ scalar}} = \frac{\kappa_1^2 v_{ew}^2}{16\pi m_{h_0}} \sqrt{1 - 4m_{\chi}^2 / m_{h_0}^2}$$

$$\Gamma_{h_0 \to \chi \chi \text{ fermion}} = \frac{A \kappa_2^2 m_{h_0}}{32\pi} \left(1 - 4m_{\chi}^2 / m_{h_0}^2\right)^{3/2}$$

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- Stronger suppression statement.

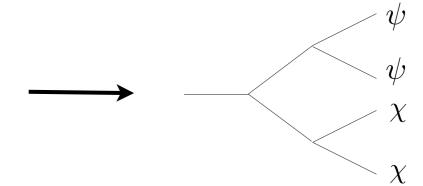
### Six Body States

(that generate two meta-stable particles + missing energy)

- Diagrams overall suppressed compared to 2, 4 body diagrams.
- Same suppression/reconstruction conclusions as the arguments for four final body states.

## Reconstructing Diagrams

Similar arguments for this diagram



Reconstruct in analogy to ZZ production.

### Meta-Stable Quarks

 Analogous charge flipping for R-hadrons interacting with the detector\*

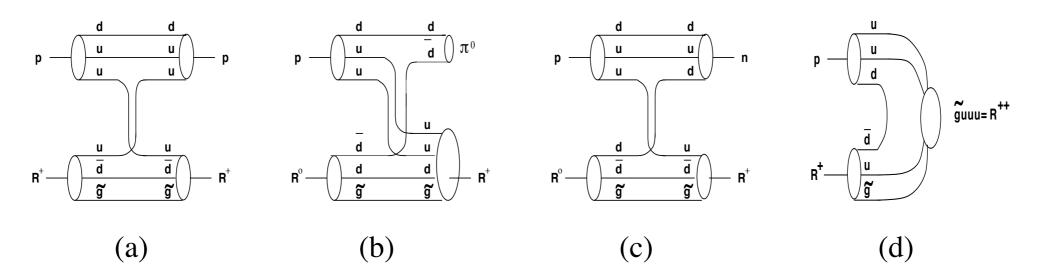


Fig. 13. *R*-hadron-proton scattering processes. (a) Elastic scattering, (b) Inelastic scattering leading to baryon and charge exchange, (c) Inelastic scattering leading to charge exchange, (d) Resonance formation.

\*From review, Fairbairn, et al. hep-ph/06 | 1040

#### Meta-Stable Particles Redux

- Particles are cleanly tagged at the LHC
  - S/B ~ O(100) for I  ${\rm fb}^{-1}$  for ATLAS and CMS.\*
- Meta-stable quarks "hadronize" into "mesons" and "baryons."
- "Mesons" and "baryons" charge flip, loose energy and even stop in the detector by exchanging valence quarks with the calorimeter.
- \* G.Aad, et al. [The ATLAS Collaboration], arXiv:0901.0512, A. Rizzi, et al. [The CMS Collaboration], CMS-AN-2007/049.

### Meta-Stable Velocities

 The "hadrons" and "leptons" look like slow, heavy "muons" in the muon detector.

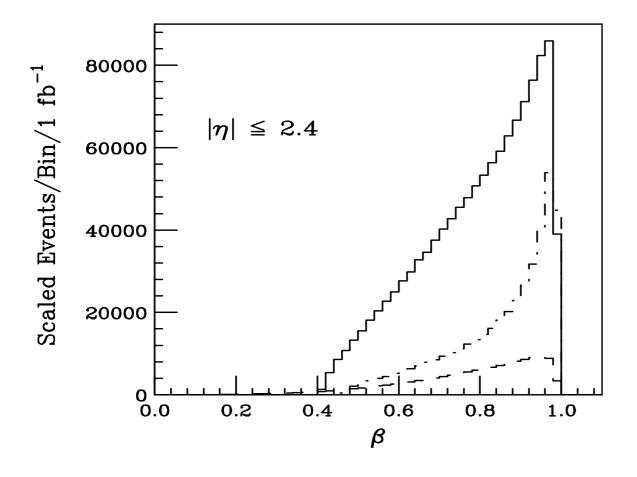


FIG. 2: Scaled events versus velocity for a 300 GeV N-hadron with (dashed) and without (solid) emission of the dark gauge boson. 200 GeV stable lepton is dot-dashed. The histograms are scaled by 20, 1 and 200, respectively.

### Time of Flight Issues

- Trigger on "heavy muons" which are out of time with rest of the relativistic event.
- "Muons" that are too slow are problematic. Can be reconstructed with the wrong event.

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- Trigger on "heavy muons" which are out of time with rest of the relativistic event.
- "Muons" that are too slow are problematic. Can be reconstructed with the wrong event.
  - LHC proton beams are collimated in bunches separated nominally by 25 nanoseconds.
  - A problem for ATLAS because of it's size. Up to three events are in the detector at one time.

#### Fast Detector Simulation

• To do the analysis wrote a fast simulation to parametrize the detector effects of Geant 3 and 4.

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 Details are not in the mainstream of this talk . . .